



Annual European Rheology Conference

co-organized with XV Meeting of the Italian Society of Rheology-SIR



Anisotropic Thermal Transport in Non-Linear Non-Isothermal Polymeric Flows

David Nieto Simavilla

Wilco M.H. Verbeeten



UNIVERSIDAD
DE BURGOS



MARIE CURIE **ACTIONS**



The MCIATTP project

A. Experimental investigation of thermal transport in polymers

- Anisotropy in thermal conductivity
- Stress-Thermal Rule
- Heat capacity vs. Deformation

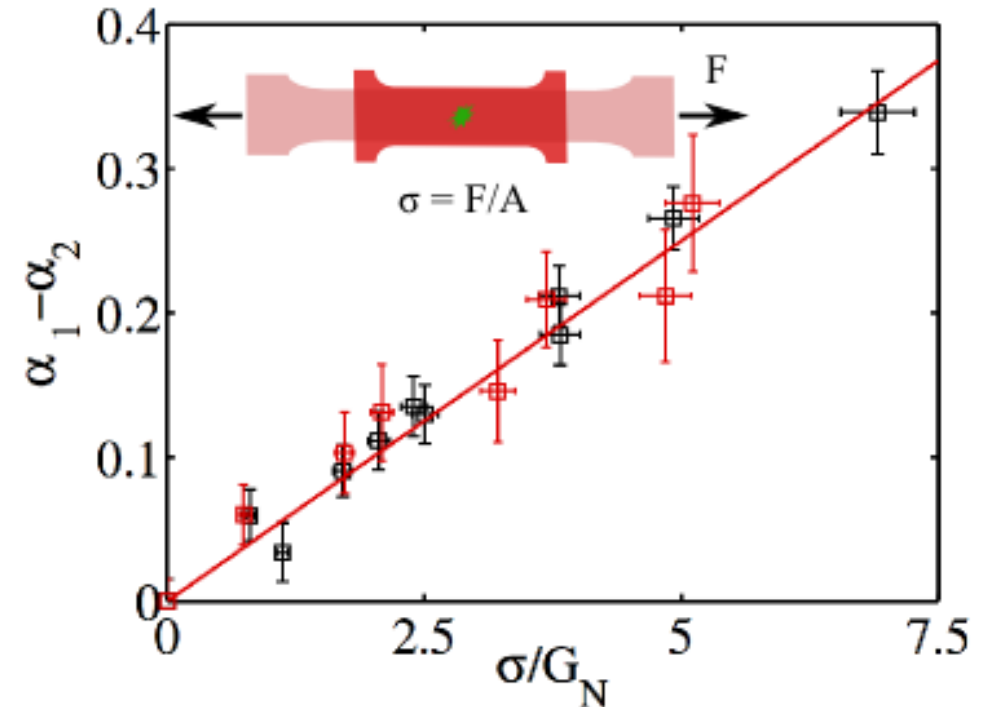
B. Implementation of constitutive models

- Branched: eXtended Pom-Pom
- Linear: Rolie Poly
- Compare predictions with available experimental data PE, PS, PMMA...

C. Develop a deeper molecular understanding

- MD Simulations
- Why universal?
- Why beyond finite extensibility?

D. Implementation of non-homogeneous non-Isothermal flow simulations



Nieto et al. J. Heat Transfer 2014

The MCIATTP project

A. Experimental investigation of thermal transport in polymers

- Anisotropy in thermal conductivity
- Stress-Thermal Rule
- Heat capacity vs. Deformation

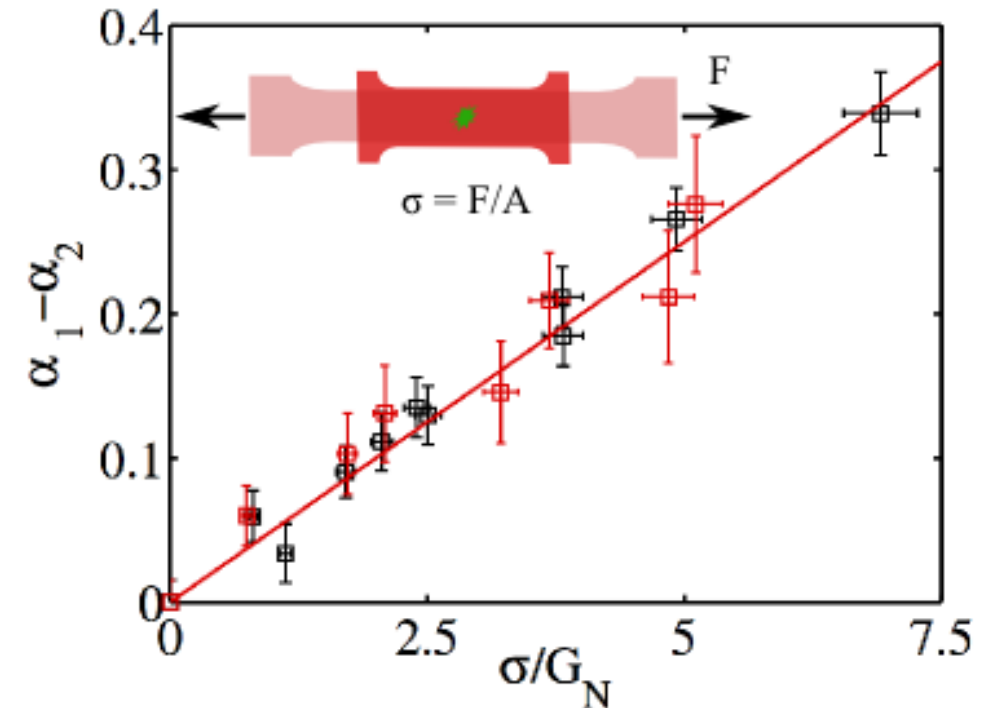
B. Implementation of constitutive models

- Branched: eXtended Pom-Pom
- Linear: Rolie Poly
- Compare predictions with available experimental data PE, PS, PMMA...

C. Develop a deeper molecular understanding

- MD Simulations
- Why universal?
- Why beyond finite extensibility?

D. Implementation of non-homogeneous non-Isothermal flow simulations



Nieto et al. J. Heat Transfer 2014

The MCIATTP project

A. Experimental investigation of thermal transport in polymers

- Anisotropy in thermal conductivity
- Stress-Thermal Rule
- Heat capacity vs. Deformation

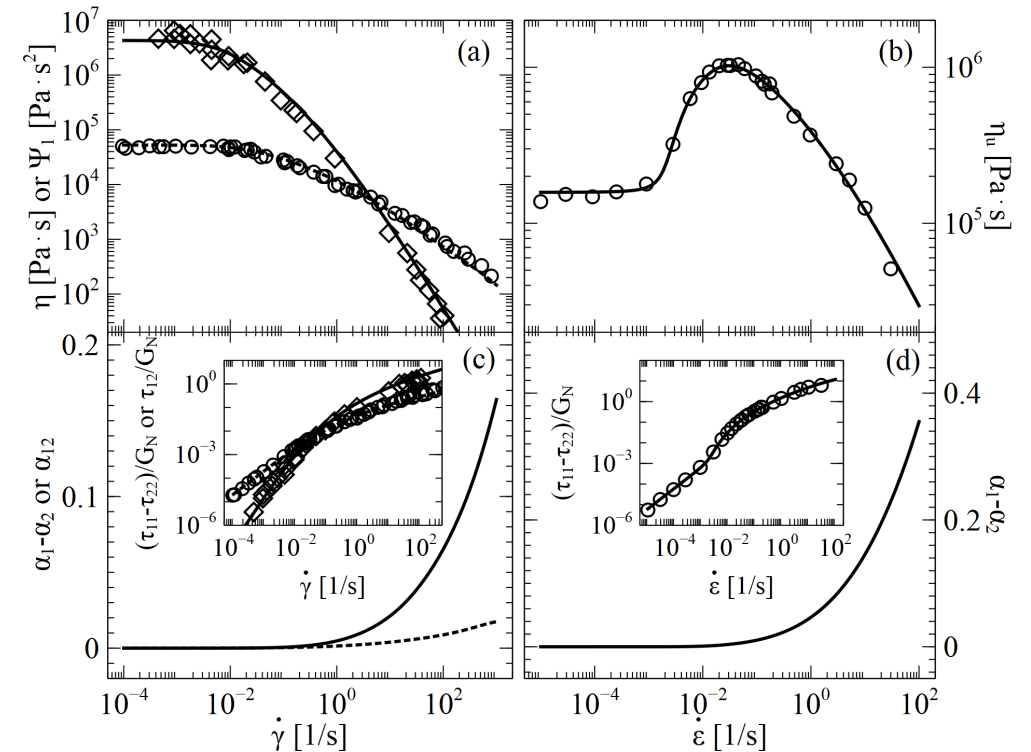
B. Implementation of constitutive models

- Branched: eXtended Pom-Pom
- Linear: Rolie Poly
- Compare predictions with available experimental data PE, PS, PMMA...

C. Develop a deeper molecular understanding

- MD simulations
- Why universal?
- Why beyond finite extensibility?

D. Implementation of non-homogeneous non-Isothermal flow simulations



The MCIATTP project

A. Experimental investigation of thermal transport in polymers

- Anisotropy in thermal conductivity
- Stress-Thermal Rule
- Heat capacity vs. Deformation

B. Implementation of constitutive models

- Branched: eXtended Pom-Pom
- Linear: Rolie Poly
- Compare predictions with available experimental data PE, PS, PMMA...

C. Develop a deeper molecular understanding

- MD Simulations
- Why universal?
- Why beyond finite extensibility?

D. Implementation of non-homogeneous non-Isothermal flow simulations



$$\mathbf{k} - \frac{1}{3}\text{tr}(\mathbf{k})\boldsymbol{\delta} \propto n_i \langle \mathbf{R}\mathbf{R} \rangle_i$$

The MCIATTP project

A. Experimental investigation of thermal transport in polymers

- Anisotropy in thermal conductivity
- Stress-Thermal Rule
- Heat capacity vs. Deformation

B. Implementation of constitutive models

- Branched: eXtended Pom-Pom
- Linear: Rolie Poly
- Compare predictions with available experimental data PE, PS, PMMA...

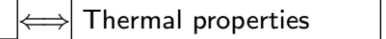
C. Develop a deeper molecular understanding

- MD Simulations
- Why universal?
- Why beyond finite extensibility?

D. Implementation of non-homogeneous non-Isothermal flow simulations

Mechanical behavior and flow

Thermal properties

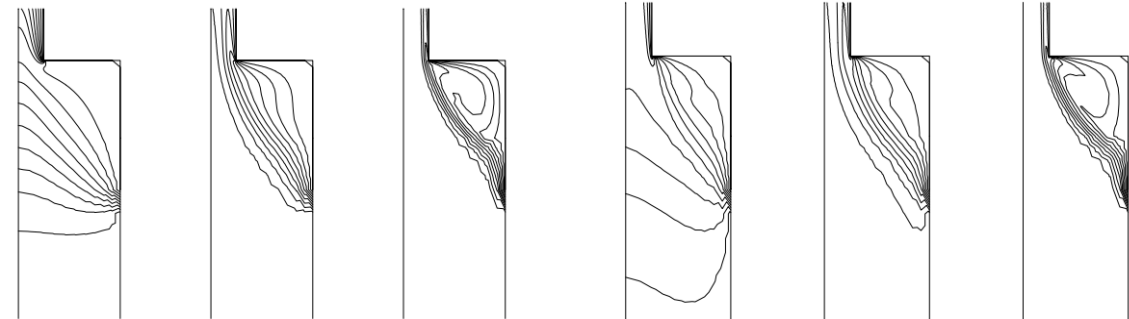


Isotropic Thermal Conductivity: k

Anisotropic Thermal Conductivity: \mathbf{k}

$$\mathbf{q} = -k\nabla T$$

$$\mathbf{q} = -\mathbf{k} \cdot \nabla T$$

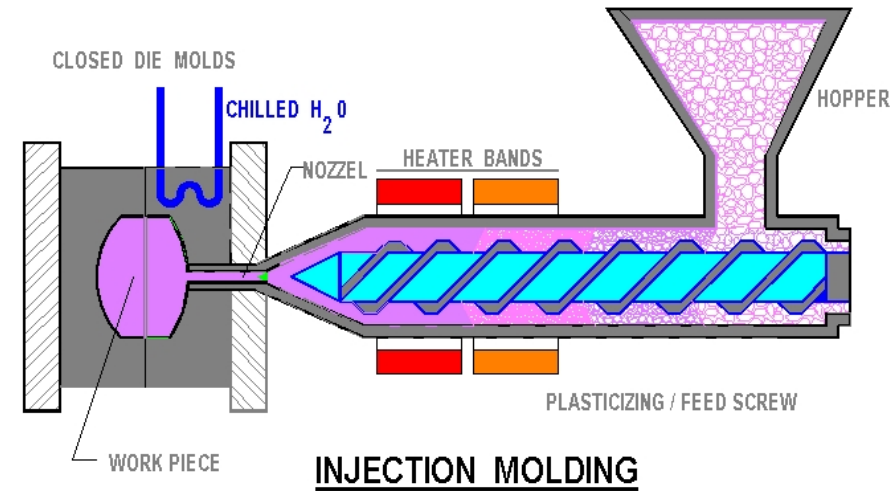


Isotherms for Pe = 10, Pe = 100 and Pe = 1000.

Isotherms for Pe = 10, Pe = 100 and Pe = 1000.

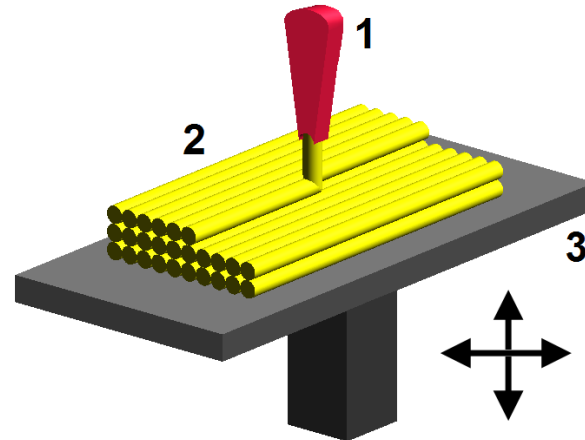
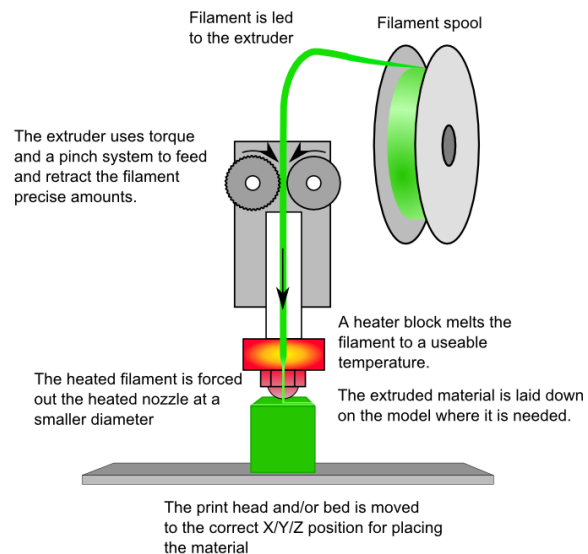
Motivation: Polymer Processing

Global plastics market is expected to reach 654 billion USD by 2020



Thermal Transport Affects:

- Injection Pressure
- Cavity Flow
- Residual Stress
- Part Shrinkage



Non-Isothermal Transport Phenomena

Balance Equations:

$$\text{Mass: } \frac{\partial \rho}{\partial t} = -\nabla \cdot (\rho \mathbf{v})$$

$$\text{Momentum: } \frac{\partial \rho \mathbf{v}}{\partial t} = -\nabla \cdot (\rho \mathbf{v} \mathbf{v} + \boldsymbol{\pi})$$

$$\text{Internal Energy: } \frac{\partial \rho \hat{u}}{\partial t} = -\nabla \cdot (\rho \hat{u} \mathbf{v} + \mathbf{q}) - \boldsymbol{\pi} : \nabla \mathbf{v}$$

Constitutive equations:

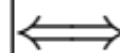
$$\mathbf{q} = -k \nabla T$$

$$\hat{c}_v = \hat{c}_v(T)$$

$$\boldsymbol{\tau} = \eta(T) [\nabla \mathbf{v} + \nabla \mathbf{v}^T]$$

- High stresses & Low thermal conductivity.

Mechanical behavior and flow



Thermal properties

Anisotropic Thermal Conduction

Fourier's Law: Thermal transport in deformed polymers is diffusive and anisotropic.

$$\mathbf{q} = -\mathbf{k} \cdot \nabla T$$

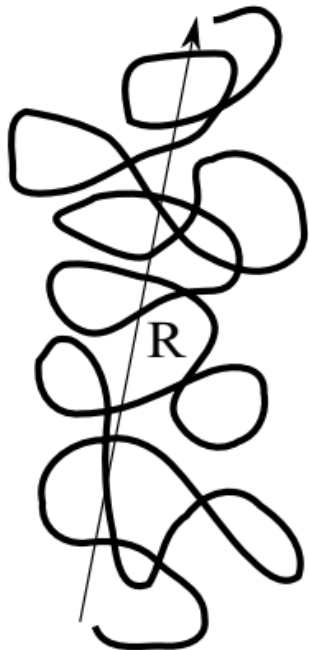
\mathbf{k} is a tensor!

Observation: k_{eq} increases with molecular weight.

Ueberreiter & Otto-Laupenmühlen, Kolloid Z. 1953

Hypothesis: *Energy transport along the backbone of a polymer chain is more efficient than between chains.*

Simple molecular arguments:



$$\mathbf{k} \propto \langle \mathbf{R}\mathbf{R} \rangle \quad + \quad \boldsymbol{\tau} \propto \langle \mathbf{R}\mathbf{R} \rangle$$

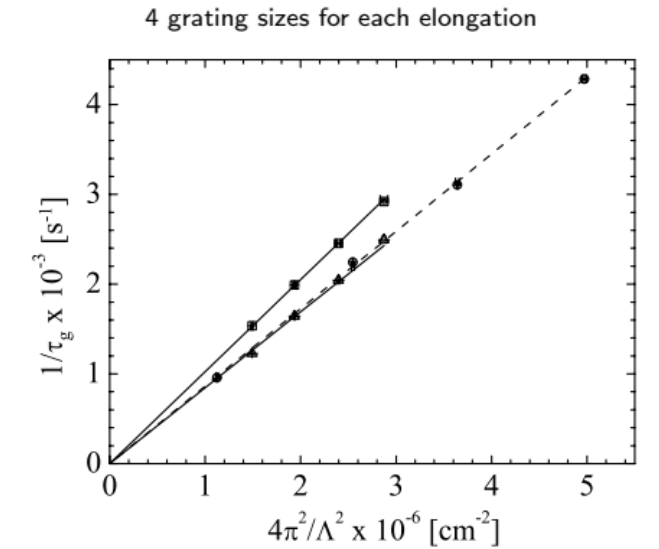
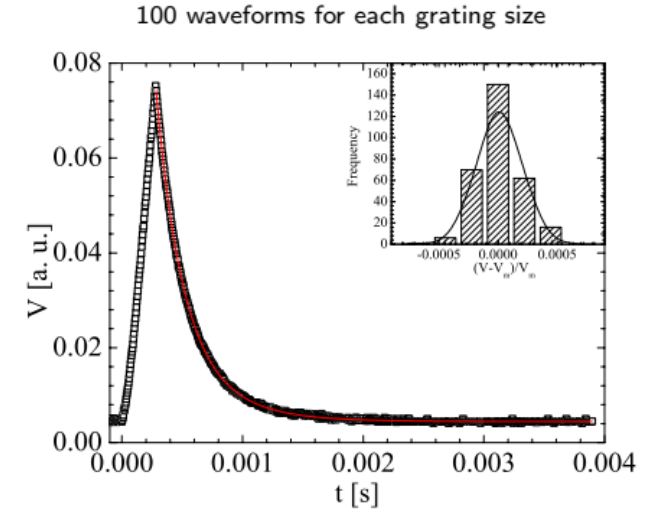
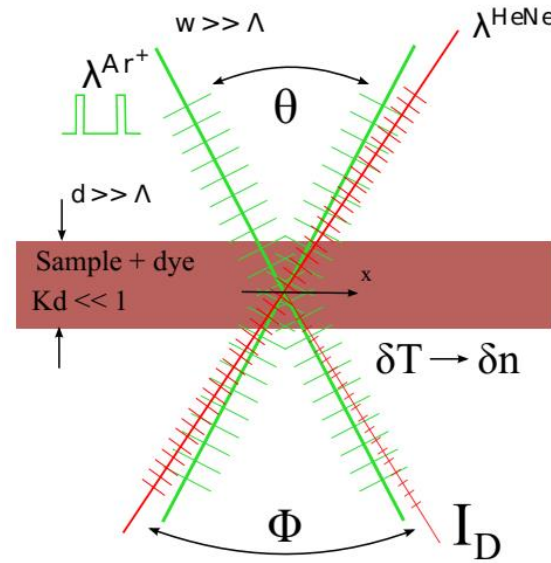
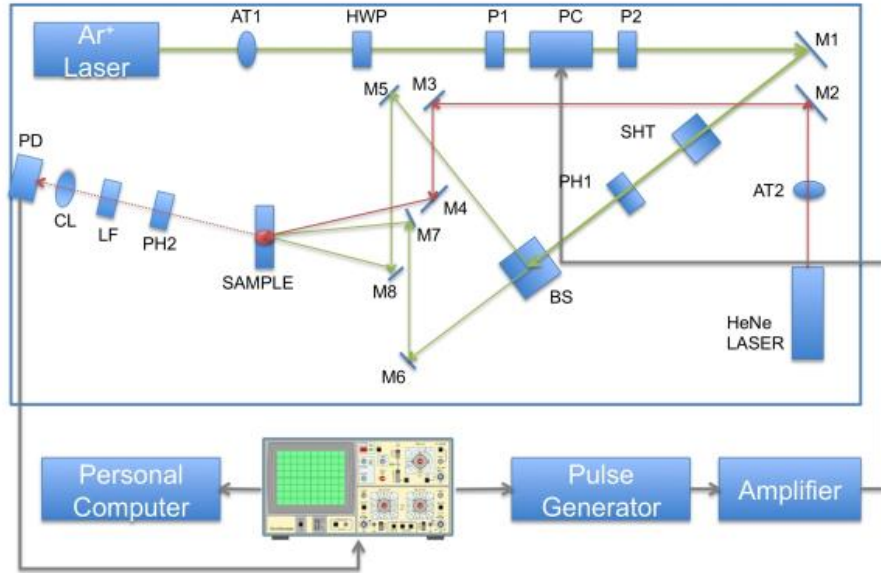
$$\mathbf{k} - \frac{1}{3}\text{tr}(\mathbf{k})\boldsymbol{\delta} = k_{\text{eq}}C_t \left[\boldsymbol{\tau} - \frac{1}{3}\text{tr}(\boldsymbol{\tau})\boldsymbol{\delta} \right]$$

The Stress-Thermal Rule

B.H.A.A. van den Brule, Rheol Acta 1989.
Öttinger and Petrillo, J. Rheol. 40 (5) 1996.
Curtiss and Bird, J. Chem. Phys. 107 (13) 1997.

$$C_t \propto \frac{nk_B^2 T}{\zeta}$$

Experiments: Forced Rayleigh Scattering (FRS)

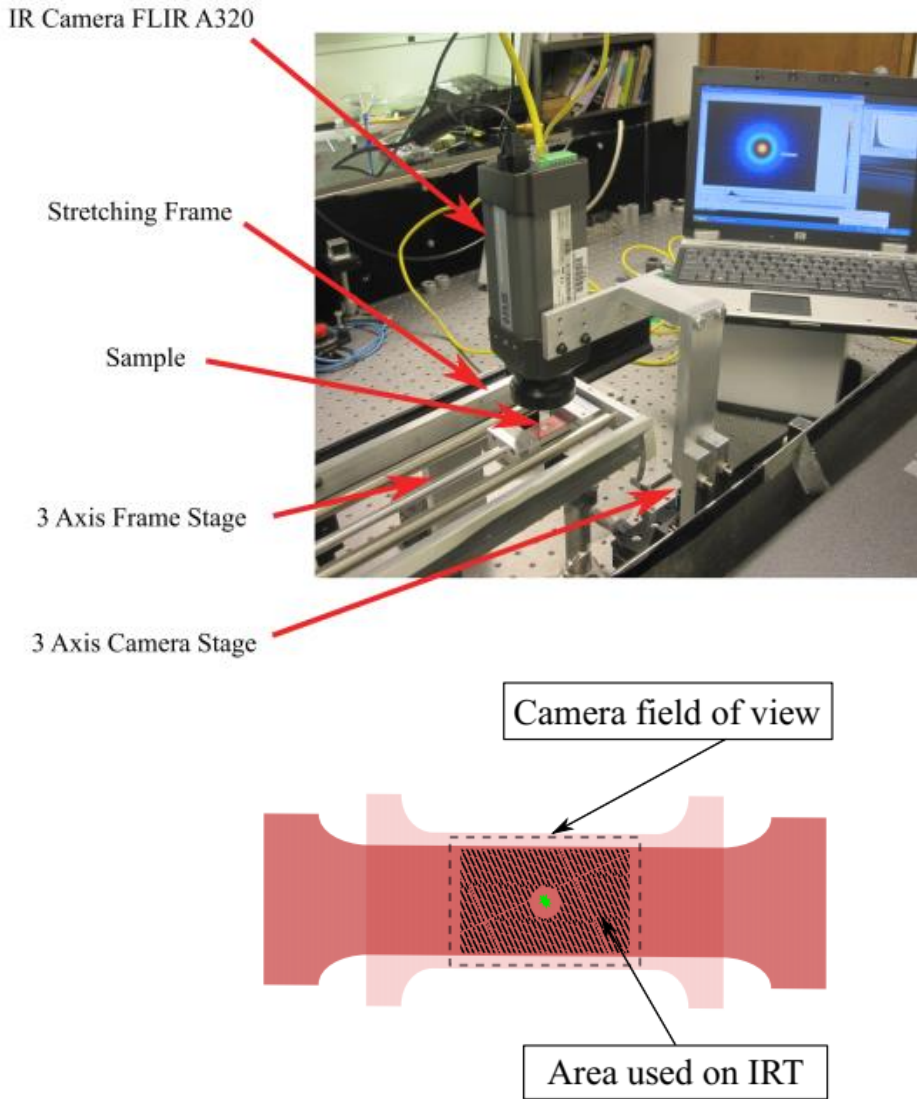


Intensity/Voltage at the photodetector:

$$V(t) = A \exp\left(-2\frac{t}{\tau_g}\right) + B \exp\left(-\frac{t}{\tau_g}\right) + C$$

$$\frac{1}{\tau_g} = D_{th} \frac{4\pi^2}{\Lambda^2} \quad D_{th} = \frac{k}{\rho \hat{c}_p}$$

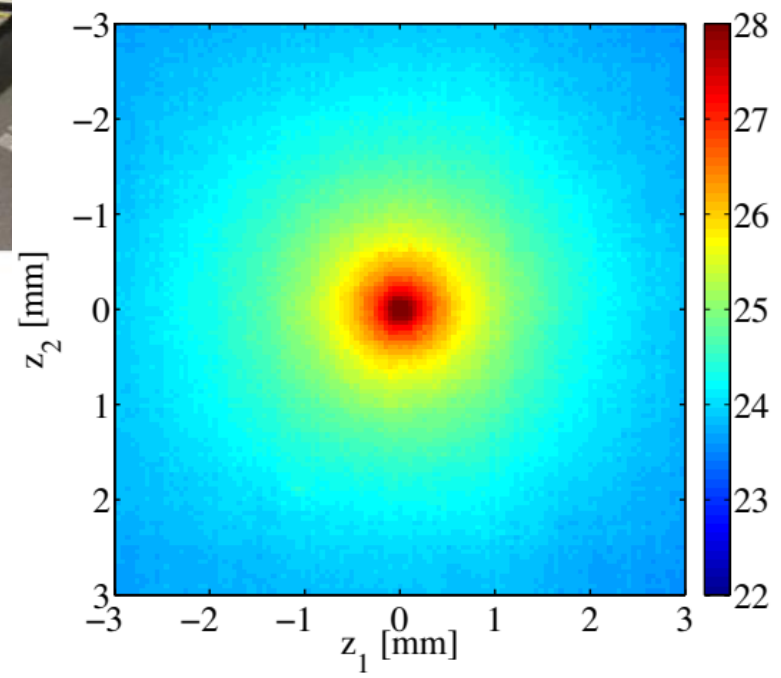
Experiments: Infrared Thermography (IRT)



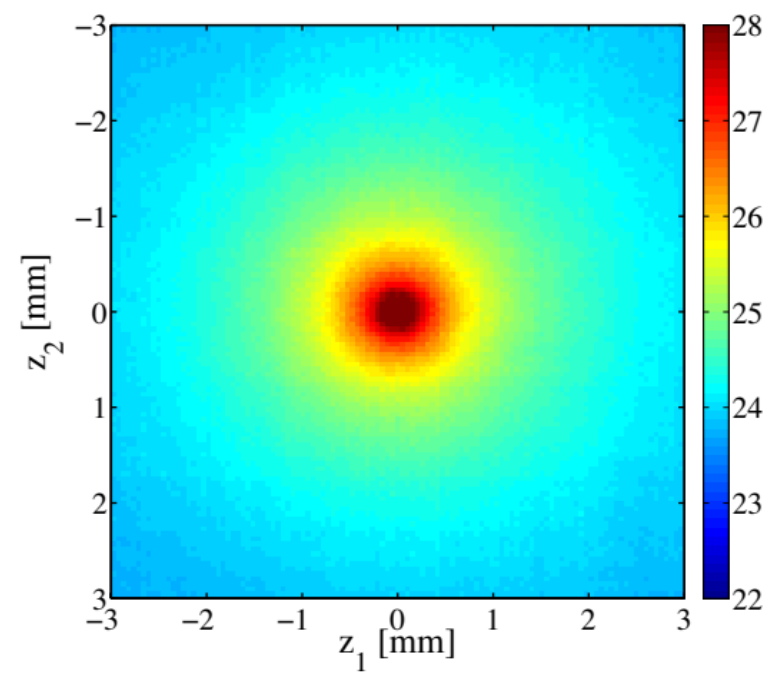
$$\theta(x_1, x_2) = \frac{1}{4\sqrt{\alpha_1\alpha_2}} K_0 \left(\sqrt{2\text{Bi}(x_1^2/\alpha_1 + x_2^2/\alpha_2)} \right)$$

$$KI_0w^2/k_{\text{eq}}, \quad \text{Bi} = hd/k_{\text{eq}}$$

$$\alpha_1 = k_{11}/k_{\text{eq}}, \quad \alpha_2 = k_{22}/k_{\text{eq}}$$

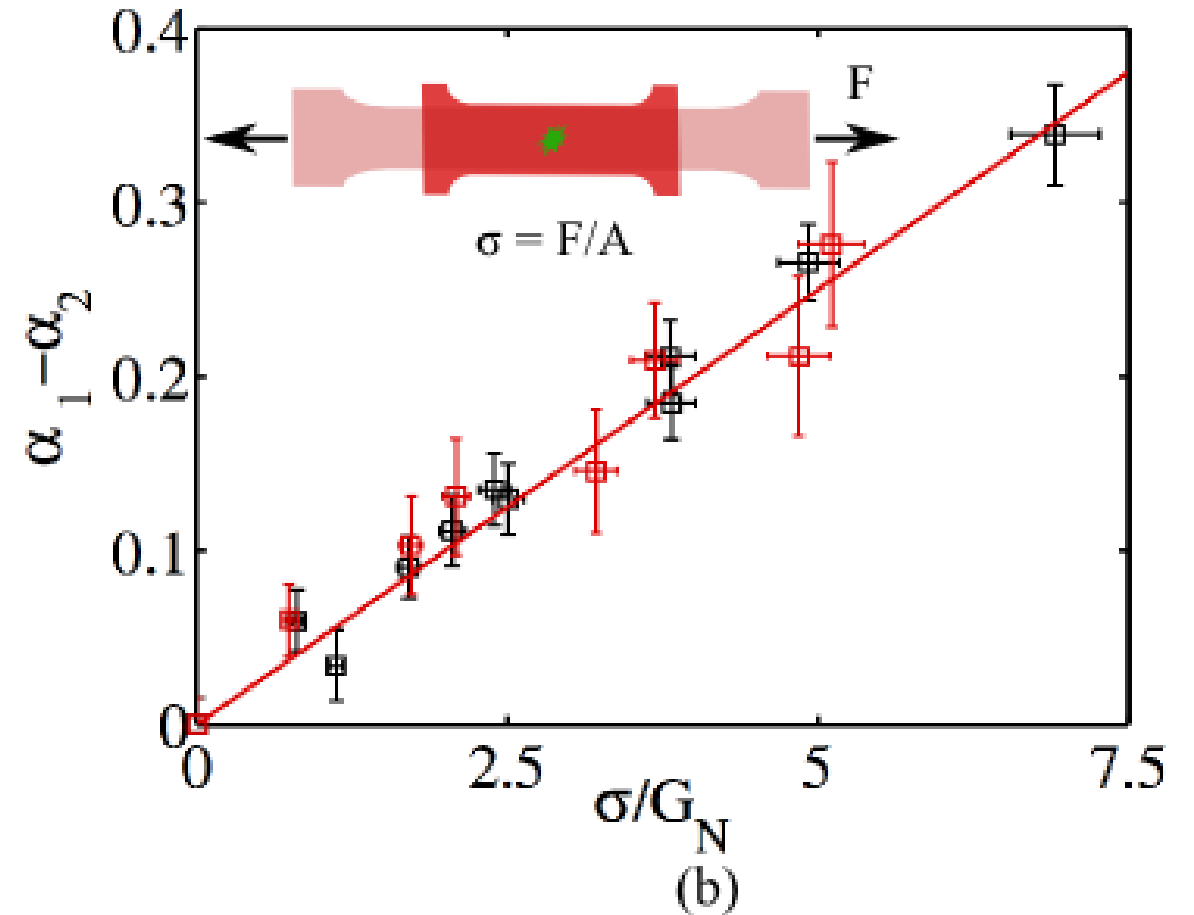
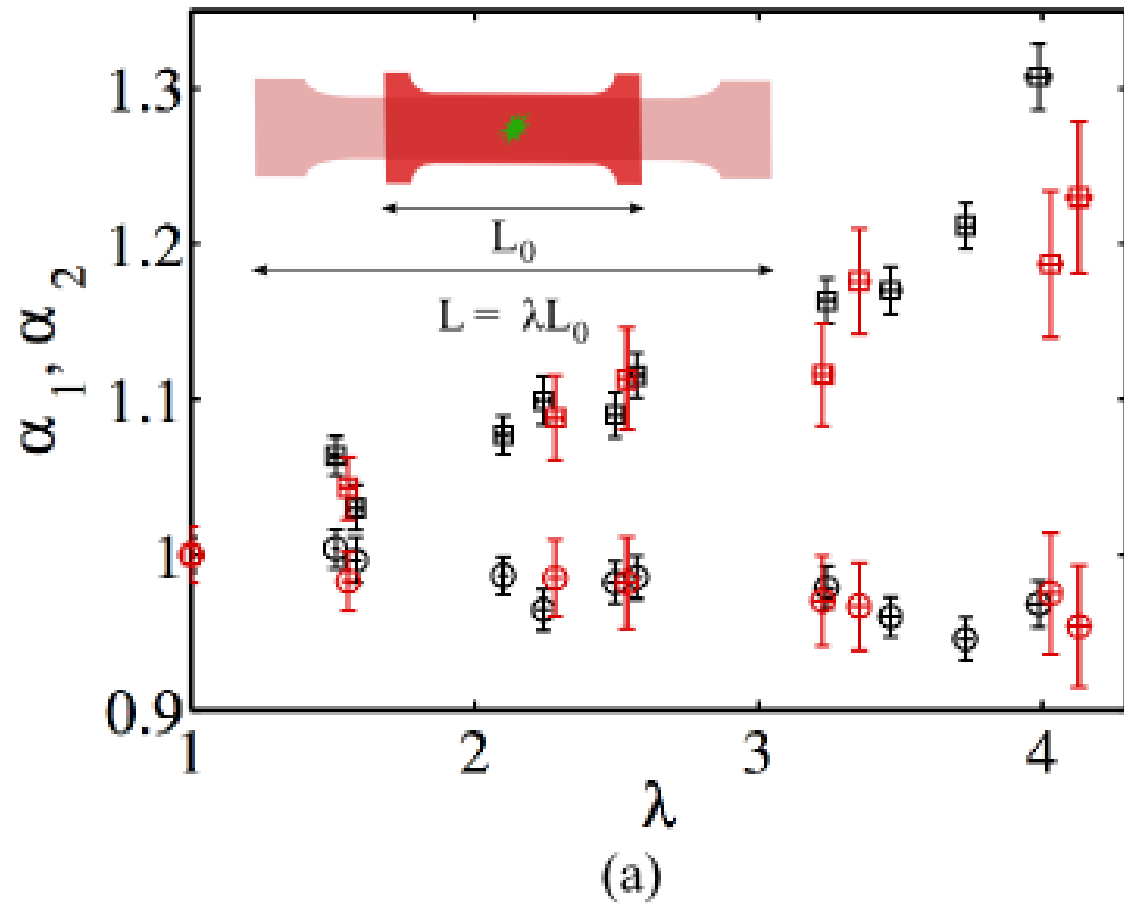


Un-stretched sample, $\lambda = 1$
and $\text{Bi}_0 = 0.029 \pm 0.001$



Stretched sample, $\lambda = 4.129$,
 $\alpha_1 = 1.23 \pm 0.049$ and $\alpha_2 = 0.954 \pm 0.039$

Comparison FRS and IRT



Key Findings: Universality...

Stress-Thermal Coefficients for several polymeric materials

Material	Deformation –	G_N [kPa]	$C_t \times 10^4$ [kPa ⁻¹]	$C_t G_N$ –	$C \times 10^9$ [Pa ⁻¹]
PIB 85k ⁷	Shear	320 ¹	1.9	0.061 ± 0.024	1.45
PIB 130k ⁷	Shear	320 ¹	1.2	0.038 ± 0.022	1.45
xI-PDMS ⁶	Uniax.	200 ¹	1.3	0.026 ± 0.008	0.13-0.26
xI-PBD 200k ⁵	Uniax.	760 ¹	0.73	0.051 ± 0.011	3.5
xI-PBD 150k ⁵	Uniax.	760 ¹	0.93	0.059 ± 0.014	3.5
xI-PI 100k ⁴	Uniax.	370 ²	0.37	0.014 ± 0.005	2.2
PS 260k ³	Uniax.	200 ¹	1.65	0.033 ± 0.007	-4.8
PMMA 83k ³	Uniax.	310 ¹	1.7	0.054 ± 0.011	0.16

$$C_t G_N \sim 0.04$$

- (1) Fetters et al. Macromolecules 27, 17 (1994)
- (2) Fetters et al. Macromolecules 37 (2004)
- (3) Gupta et al. Journal of Rheology 57 (2013)
- (4) Nieto Simavilla et al. J. Pol. Sci. B 50 (2012)
- (5) Venerus et al. Macromolecules 42 (2009)
- (6) Broerman et al. J.Chem. Phys. 111 (1999)
- (7) Venerus et al. Phys. Rev. Lett. 82 (1999)

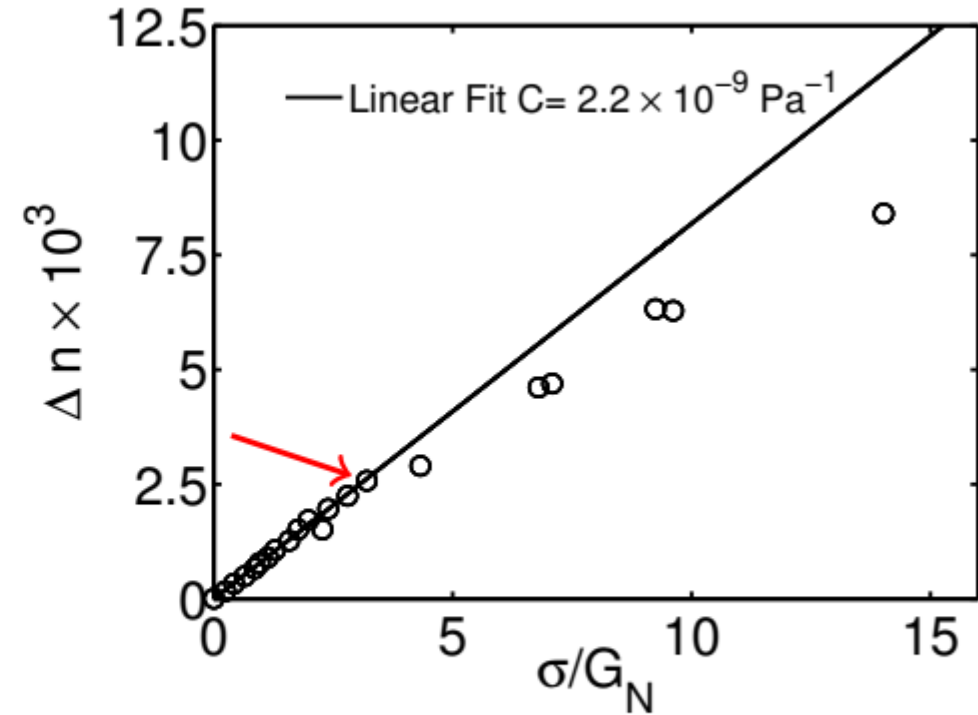
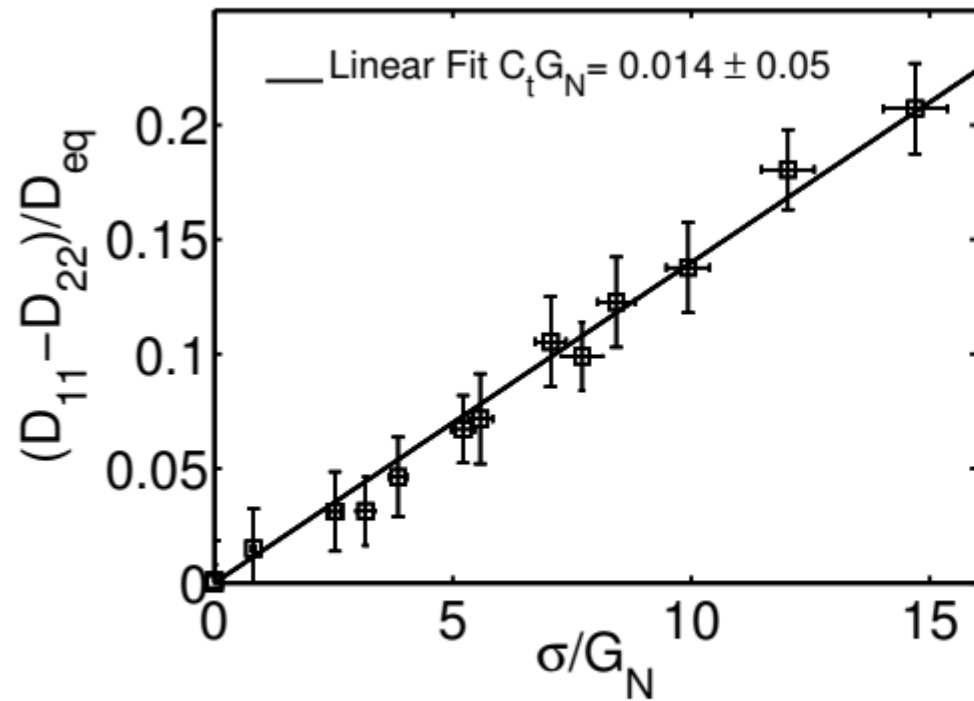
Stress-thermal Rule:

$$\mathbf{k} - \frac{1}{3} \text{tr}(\mathbf{k}) \boldsymbol{\delta} = k_{\text{eq}} C_t (\boldsymbol{\tau} - \frac{1}{3} \text{tr}(\boldsymbol{\tau}) \boldsymbol{\delta})$$

Stress-optic Rule:

$$\mathbf{n} - \frac{1}{3} \text{tr}(\mathbf{n}) \boldsymbol{\delta} = C (\boldsymbol{\tau} - \frac{1}{3} \text{tr}(\boldsymbol{\tau}) \boldsymbol{\delta})$$

Key Findings: ...Beyond Finite Extensibility



The STR stays valid where the SOR fails!

Anisotropic Thermal Conduction

Fourier's Law: Thermal transport in deformed polymers is diffusive and anisotropic.

$$\mathbf{q} = -\mathbf{k} \cdot \nabla T$$

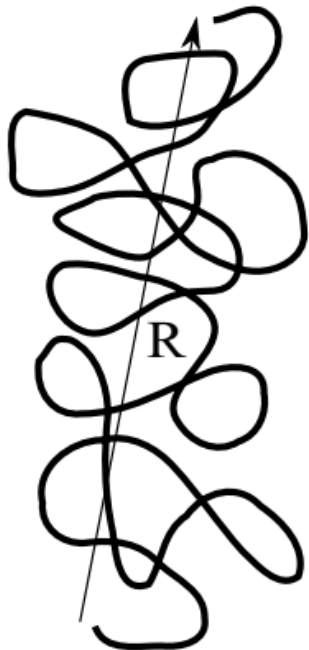
\mathbf{k} is a tensor!

Observation: k_{eq} increases with molecular weight.

Ueberreiter & Otto-Laupenmühlen, Kolloid Z. 1953

Hypothesis: *Energy transport along the backbone of a polymer chain is more efficient than between chains.*

Simple molecular arguments:



$$\mathbf{k} \propto \langle \mathbf{R}\mathbf{R} \rangle \quad + \quad \boldsymbol{\tau} \propto \langle \mathbf{R}\mathbf{R} \rangle$$

$$\mathbf{k} - \frac{1}{3}\text{tr}(\mathbf{k})\boldsymbol{\delta} = k_{\text{eq}}C_t \left[\boldsymbol{\tau} - \frac{1}{3}\text{tr}(\boldsymbol{\tau})\boldsymbol{\delta} \right]$$

The Stress-Thermal Rule

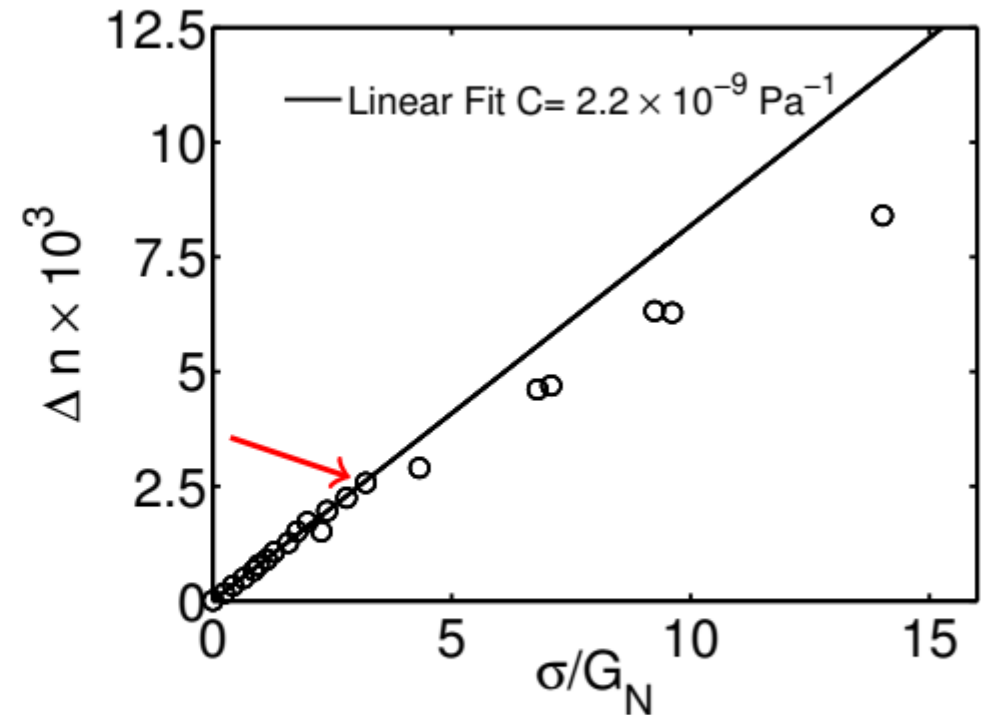
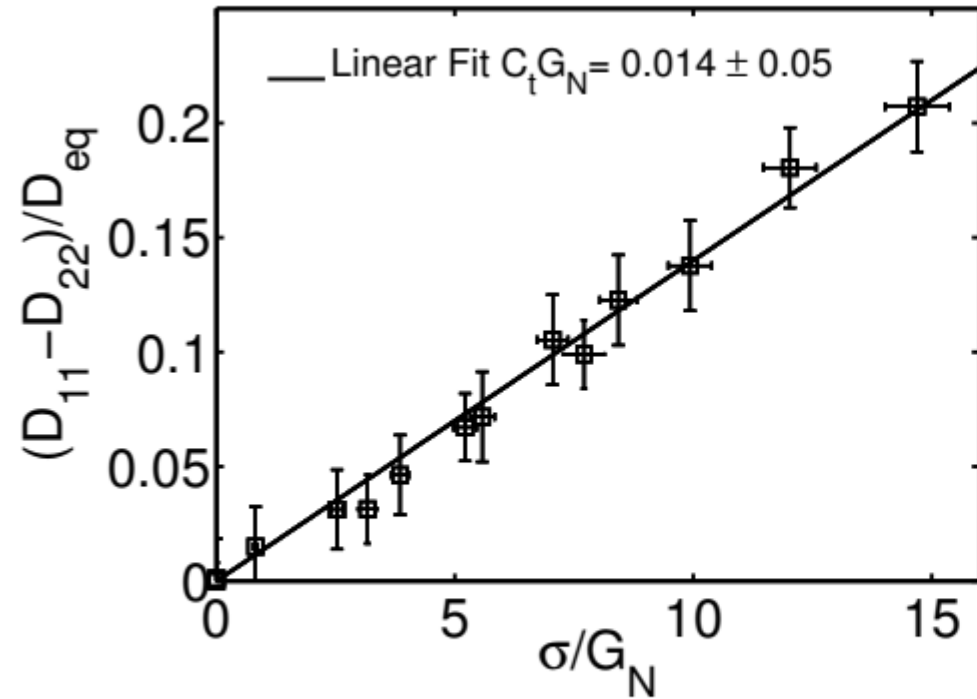
B.H.A.A. van den Brule, Rheol Acta 1989.

Öttinger and Petrillo, J. Rheol. 40 (5) 1996.

Curtiss and Bird, J. Chem. Phys. 107 (13) 1997.

$$C_t \propto \frac{nk_B^2 T}{\zeta}$$

Key Findings: ...Beyond Finite Extensibility



The STR stays valid where the SOR fails!

Nieto Simavilla et al. J. Pol. Sci. B 2012

The Stress-Thermal Rule can be applied:

1. To any melt just by knowing stress and G_N
2. At high strain and strain rates beyond the onset of finite extensibility effects

Constitutive Model: eXtended Pom-Pom

- What physics are in the model?

$$\overset{\nabla}{\boldsymbol{\tau}} + \boldsymbol{\lambda}(\boldsymbol{\tau})^{-1} \cdot \boldsymbol{\tau} - 2G_0 \mathbf{D}_u = \mathbf{0}$$

$$\alpha \neq 0 \rightarrow \Psi_2 \neq 0$$

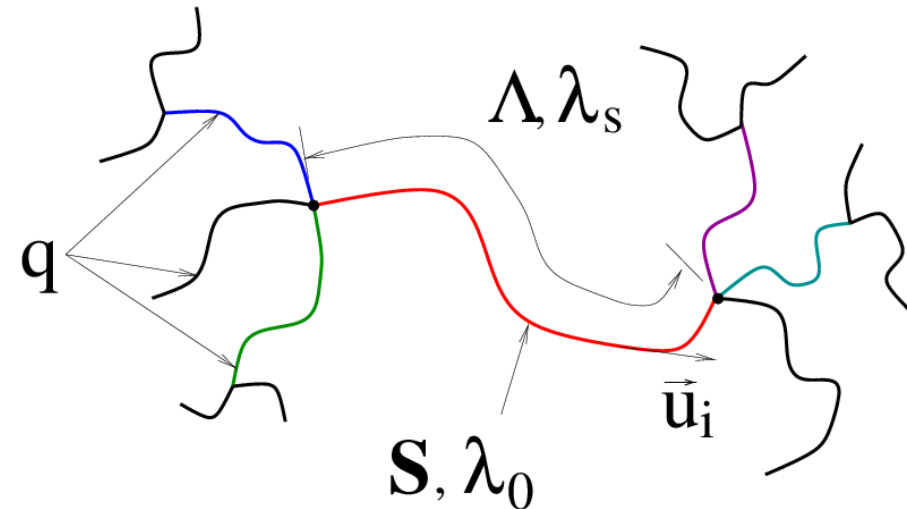
$$\boldsymbol{\lambda}(\boldsymbol{\tau})^{-1} = \frac{1}{\lambda_{0b}} \left[\frac{\alpha}{G_0} \boldsymbol{\tau} + f(\boldsymbol{\tau})^{-1} \mathbf{I} + G_0 (f(\boldsymbol{\tau})^{-1} - 1) \boldsymbol{\tau}^{-1} \right] \quad \Lambda = \sqrt{1 + \frac{I_{\boldsymbol{\tau}}}{3G_0}}$$

$$\frac{1}{\lambda_{0b}} f(\boldsymbol{\tau})^{-1} = \frac{2}{\lambda_s} \left(1 - \frac{1}{\Lambda}\right) + \frac{1}{\lambda_{0b}} \left(\frac{1}{\Lambda^2} - \frac{\alpha I_{\boldsymbol{\tau} \cdot \boldsymbol{\tau}}}{3G_0^2 \Lambda^2} \right)$$

$$\lambda_s = \lambda_{0s} e^{-\frac{2}{q}(\Lambda-1)}$$

- Why XPP?

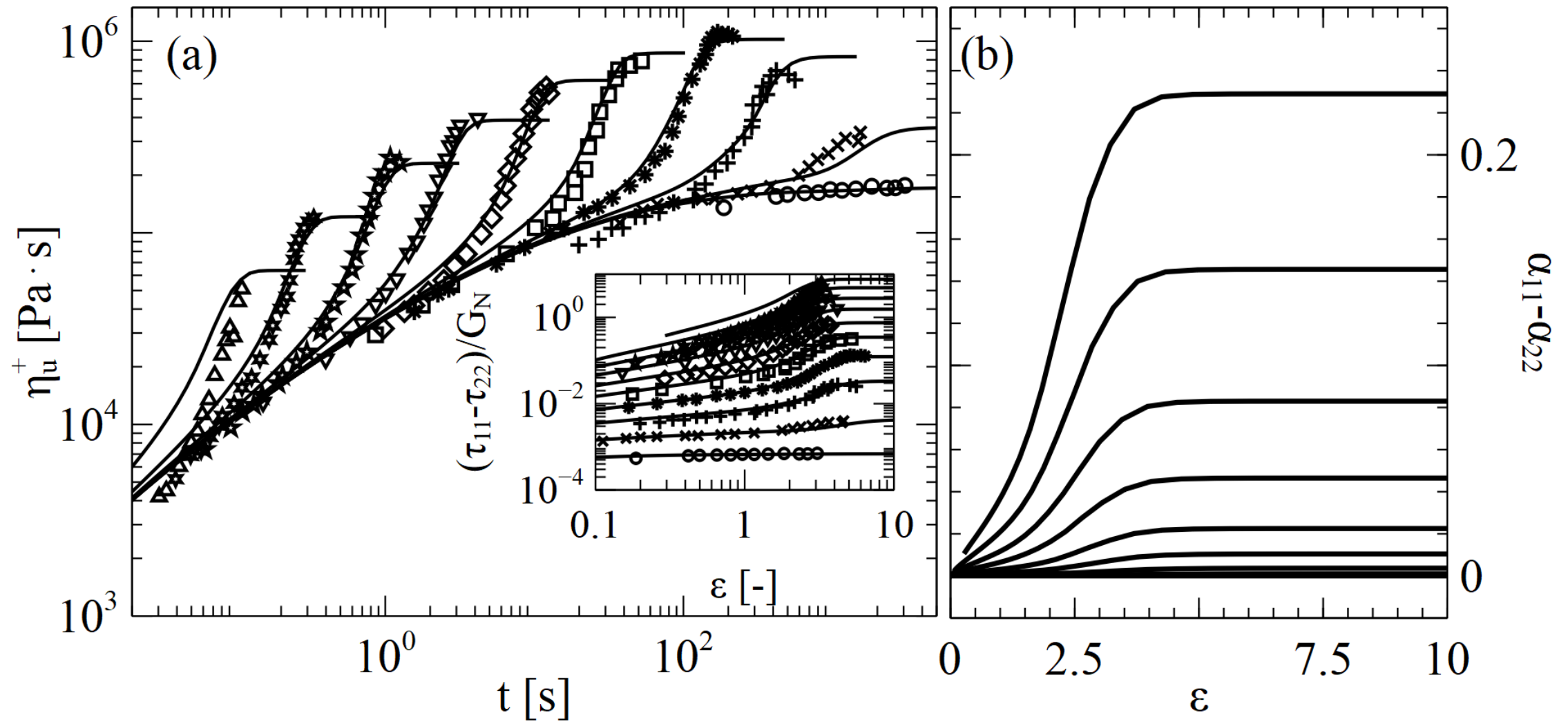
- Amenable to FEM
- Able to describe non-linear rheology
- X: Avoids finite extensibility discontinuities
- X: Includes second normal stress difference



Data: IUPAC_A LDPE melt at 170°C
Verbeeten et al. JOR 2001

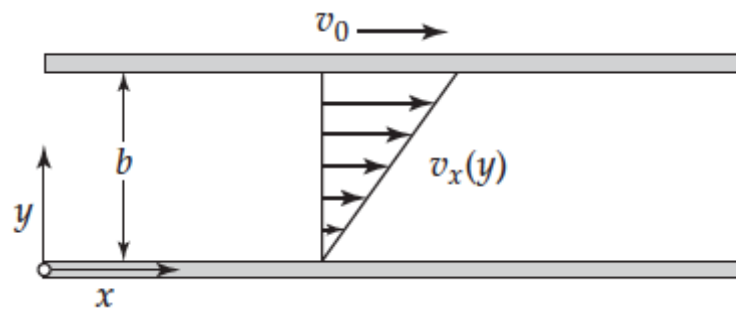
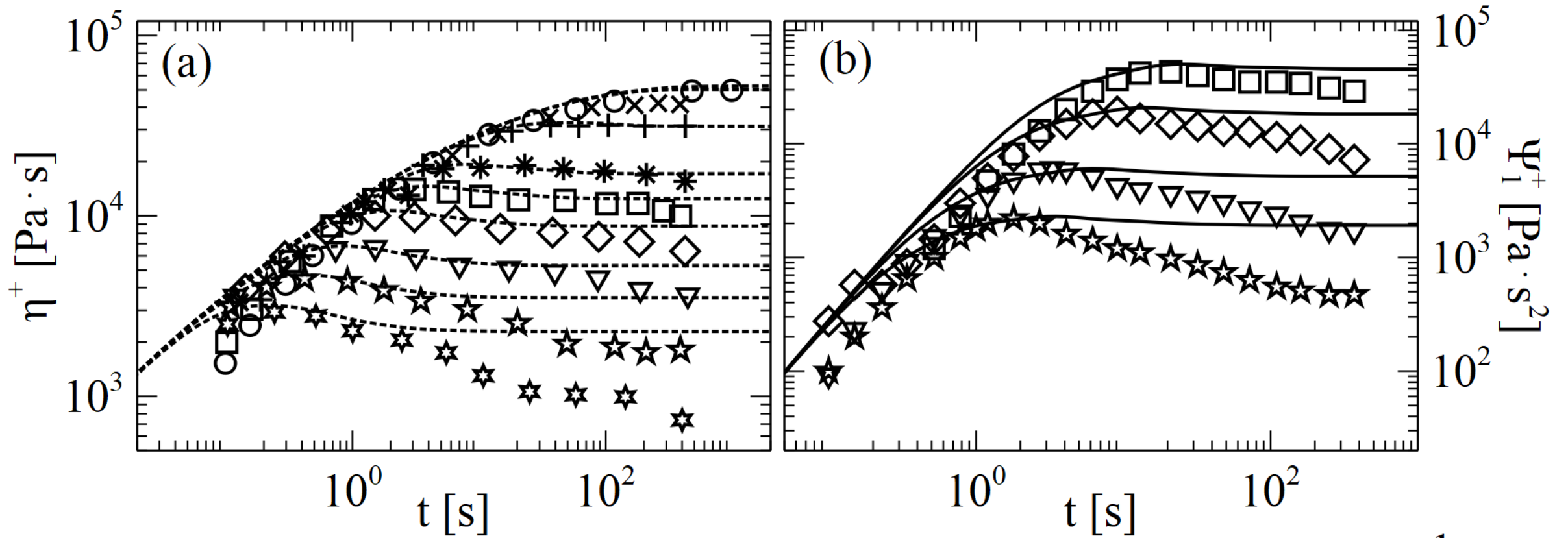
PP: McLeish and Larson. JOR 1998
xPP: Verbeeten et al. JOR 2001

Transient Start-up: Uniaxial IUPAC_A LDPE

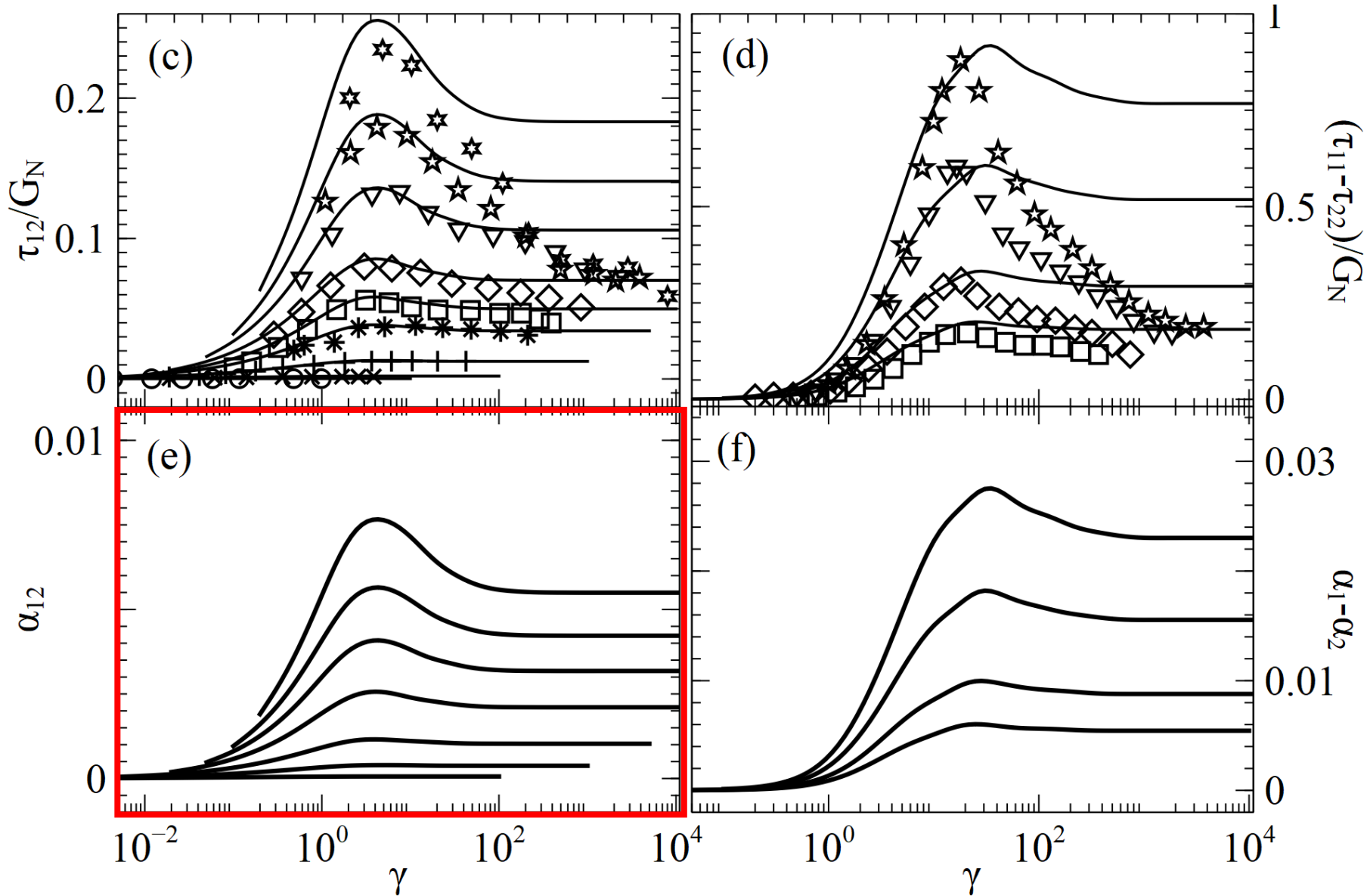


The anisotropy in TC is comparable to that observed in PS and PMMA melts $\sim 20\%$.
Gupta et al. Journal of Rheology 57, 2013.

Transient Start-up: Shear Rheology IUPAC_A LDPE



Transient Start-up: Shear IUPAC_A LDPE



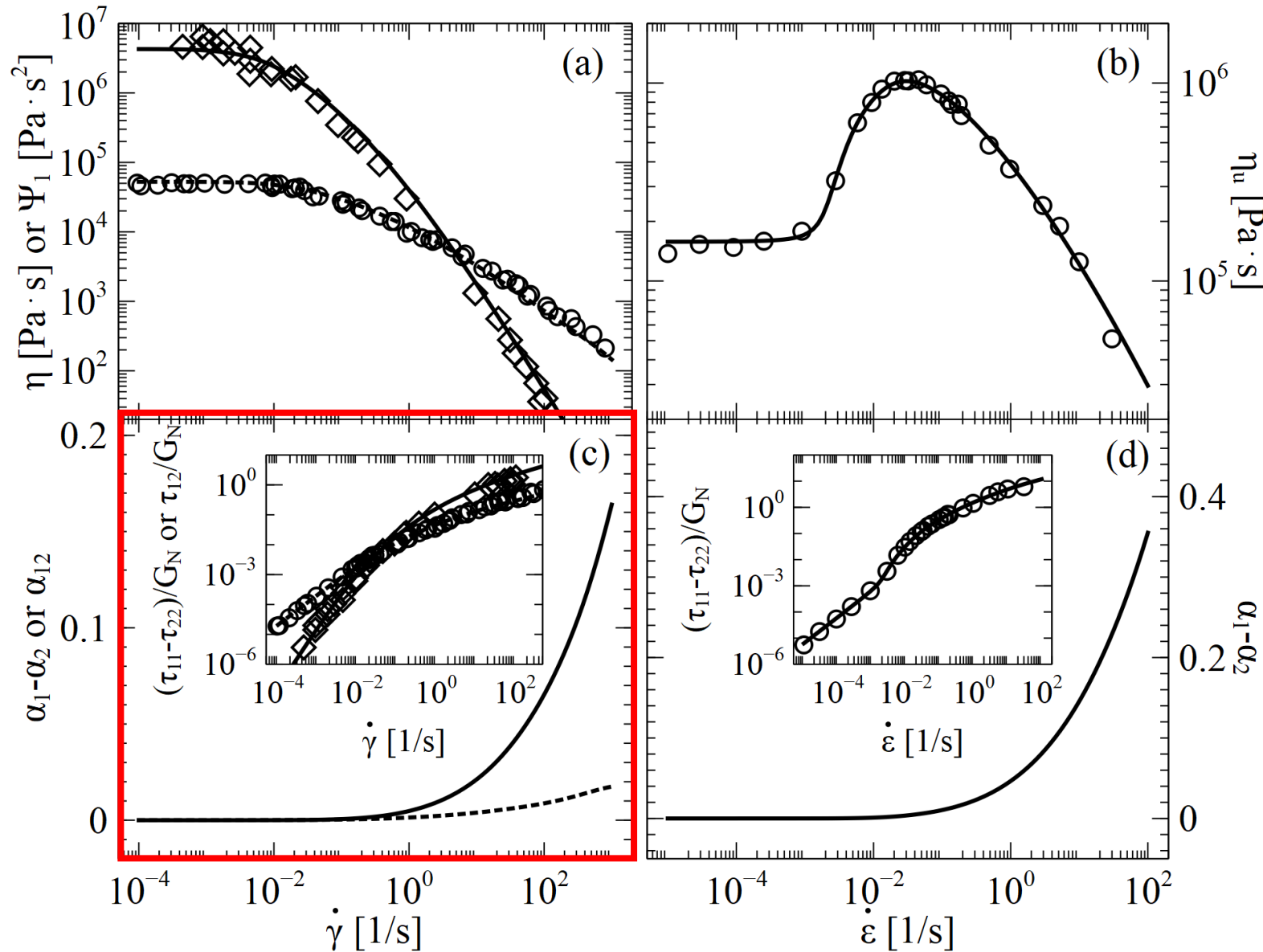
There is a non-zero off-diagonal component in shear flows

$$\mathbf{k} = \begin{pmatrix} k_{11} & k_{12} & k_{13} \\ & k_{22} & k_{23} \\ & & k_{33} \end{pmatrix}$$

$$\mathbf{q} = \mathbf{k} \cdot \nabla T$$

A temperature gradient in the 1-direction can generate heat flow in the 2-direction:
Thermal Hall Effect

Steady-State: Shear and Uniaxial Ext. IUPAC_A LDPE



There is a non-zero off-diagonal component in shear flows

$$\mathbf{k} = \begin{pmatrix} k_{11} & k_{12} & k_{13} \\ & k_{22} & k_{23} \\ & & k_{33} \end{pmatrix}$$

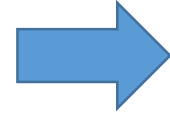
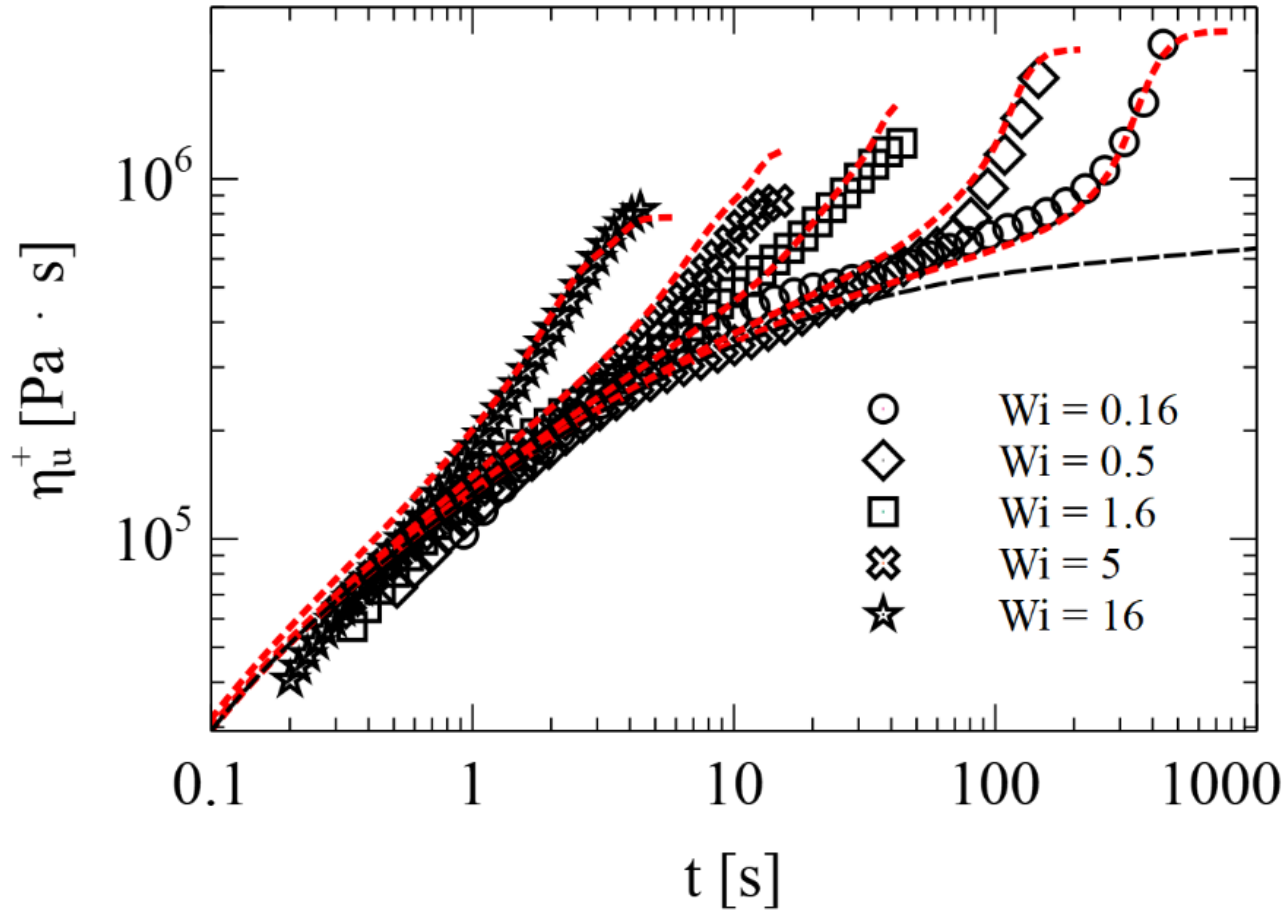
$$\mathbf{q} = \mathbf{k} \cdot \nabla T$$

A temperature gradient in the 1-direction can generate heat flow in the 2-direction:

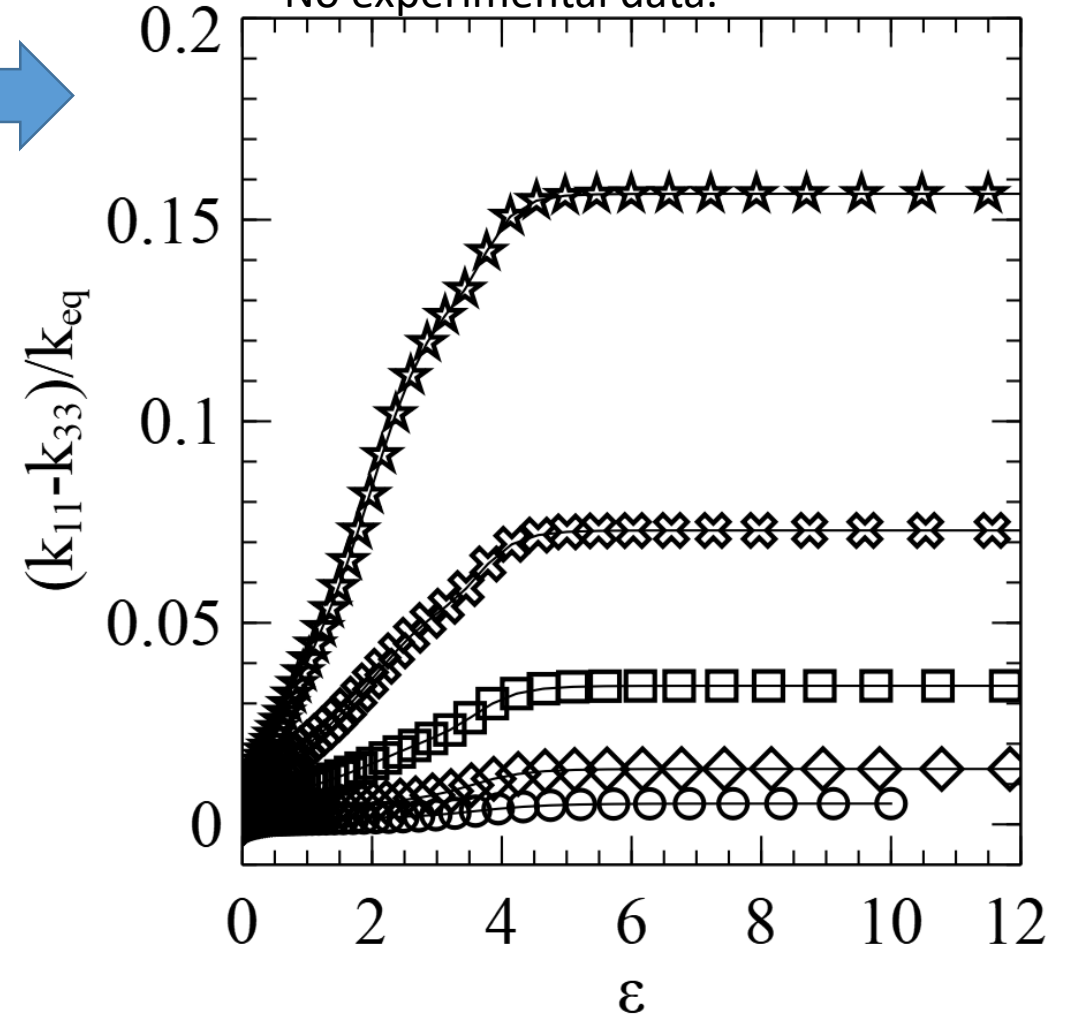
Thermal Hall Effect

Transient Start-up: Uniaxial PS158k

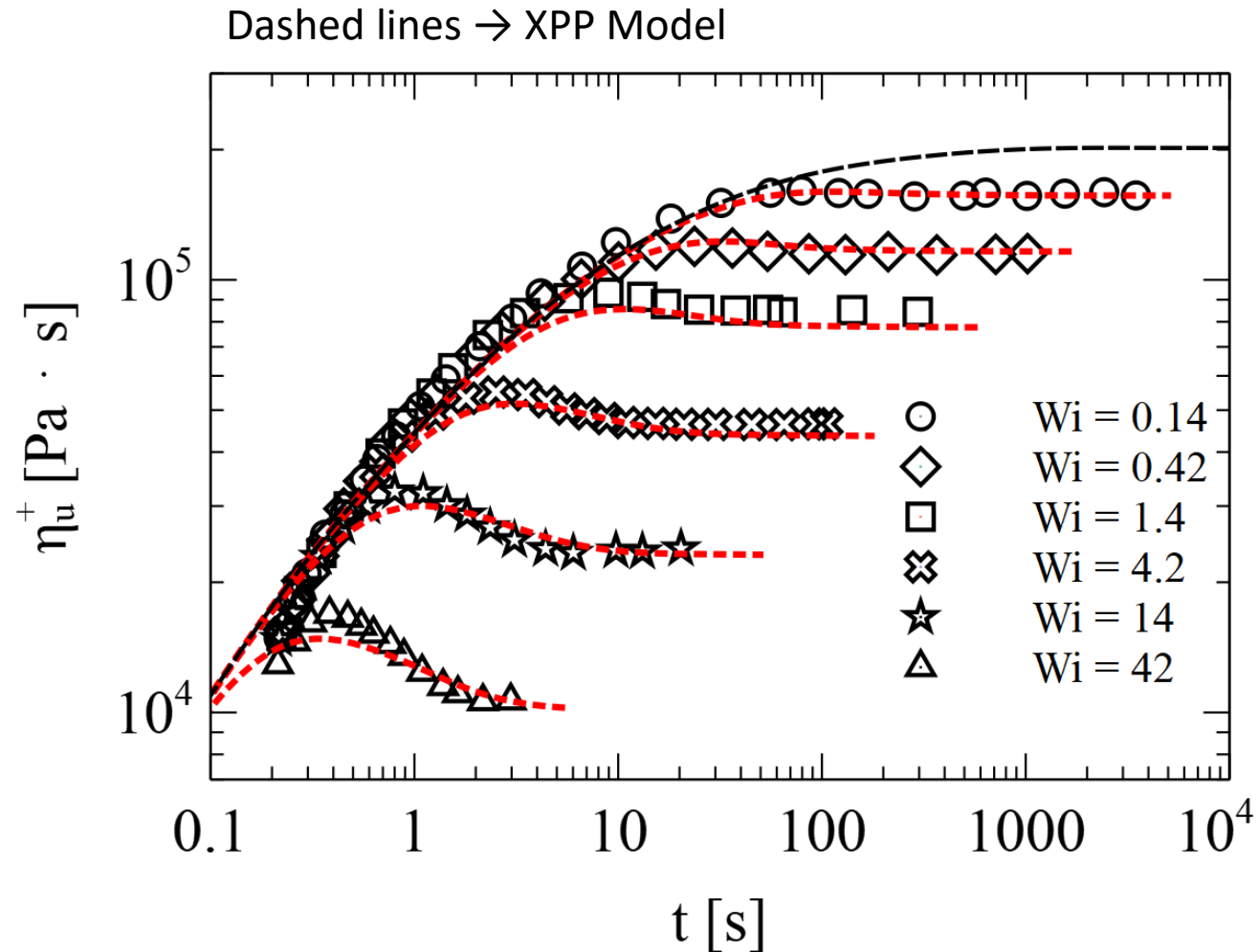
Dashed lines → XPP Model



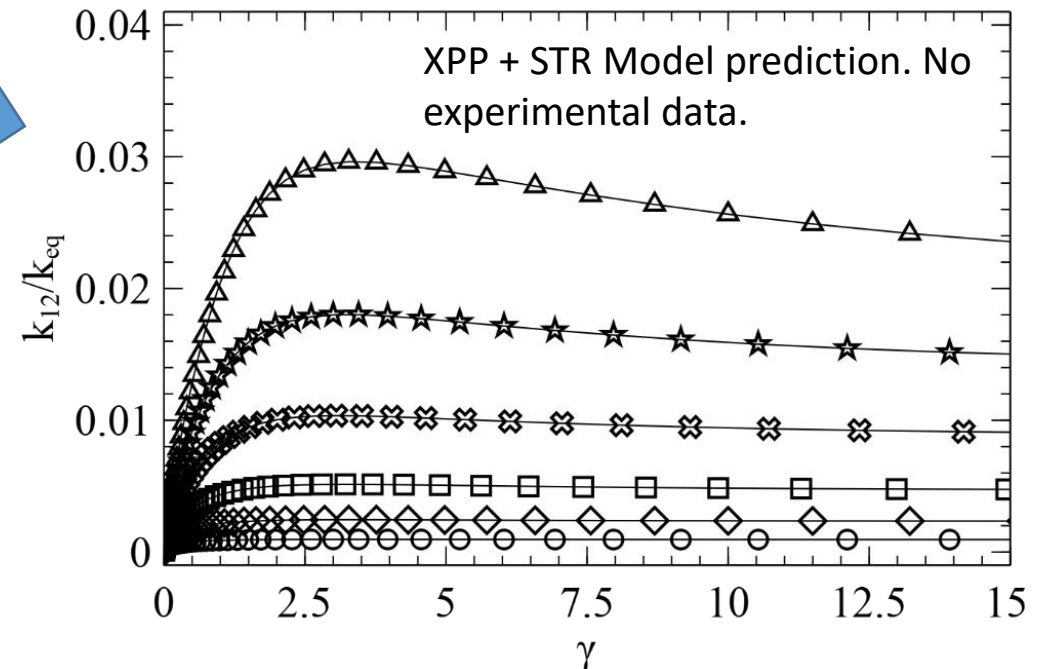
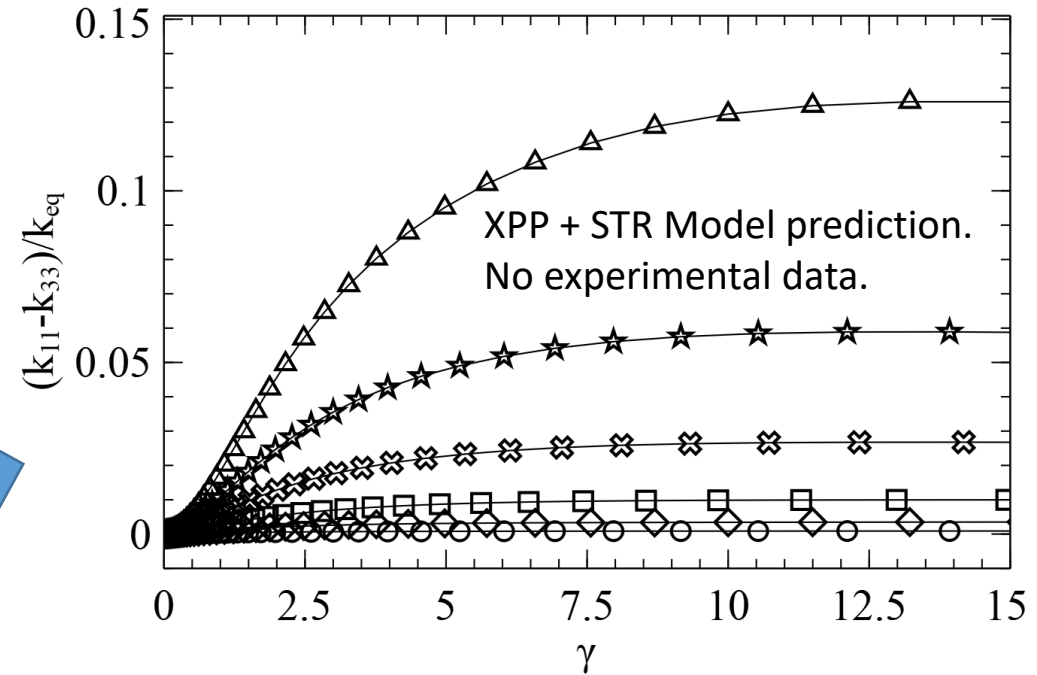
XPP + STR Model prediction.
No experimental data.



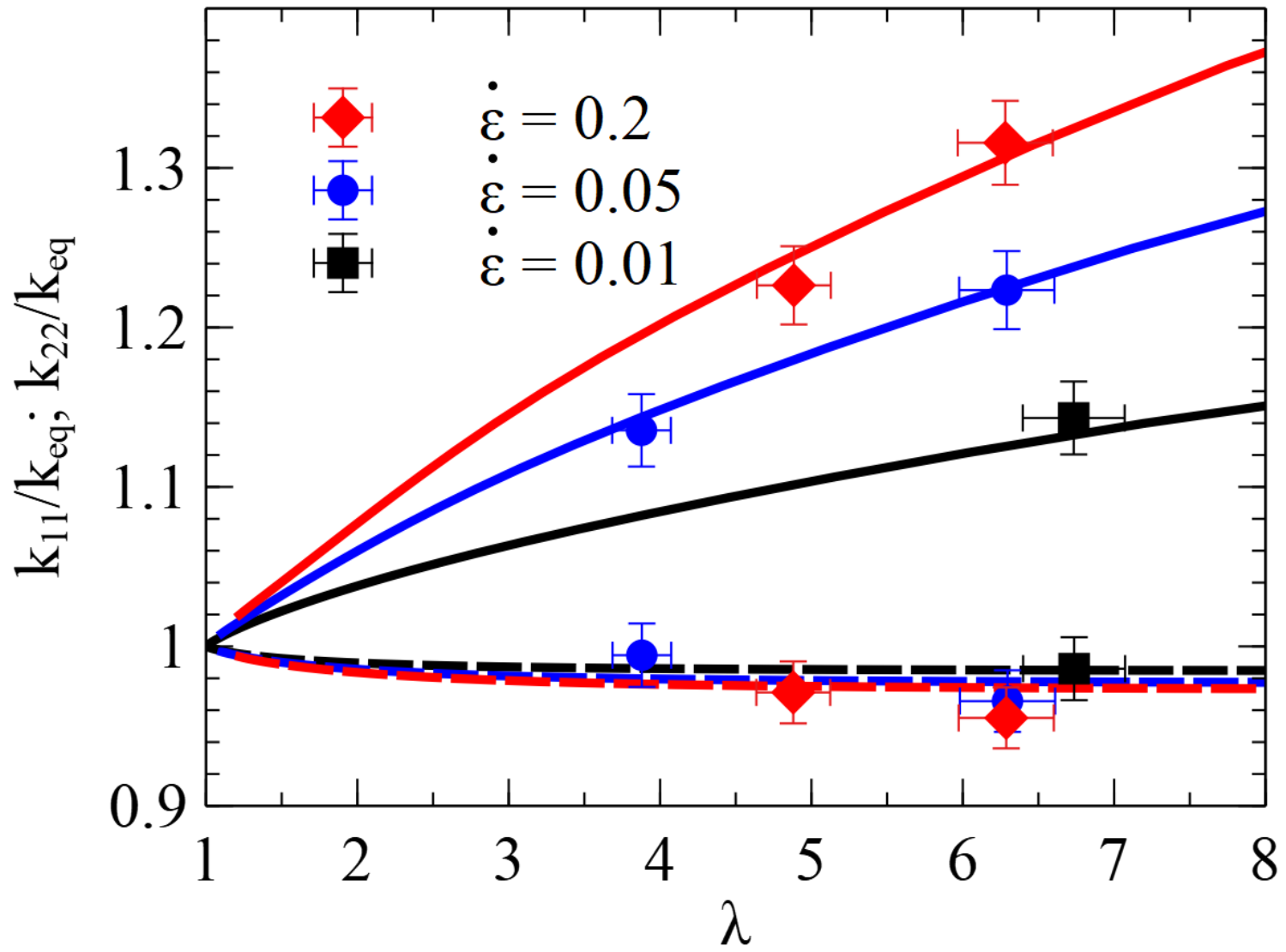
Transient Start-up: Shear Rheology PS158k



Data: Thomas Schweizer Rheol. Acta 2002



Comparison to experiments: PS260k



FRS Measurements after quenching. Data from Gupta et al. JOR 2013

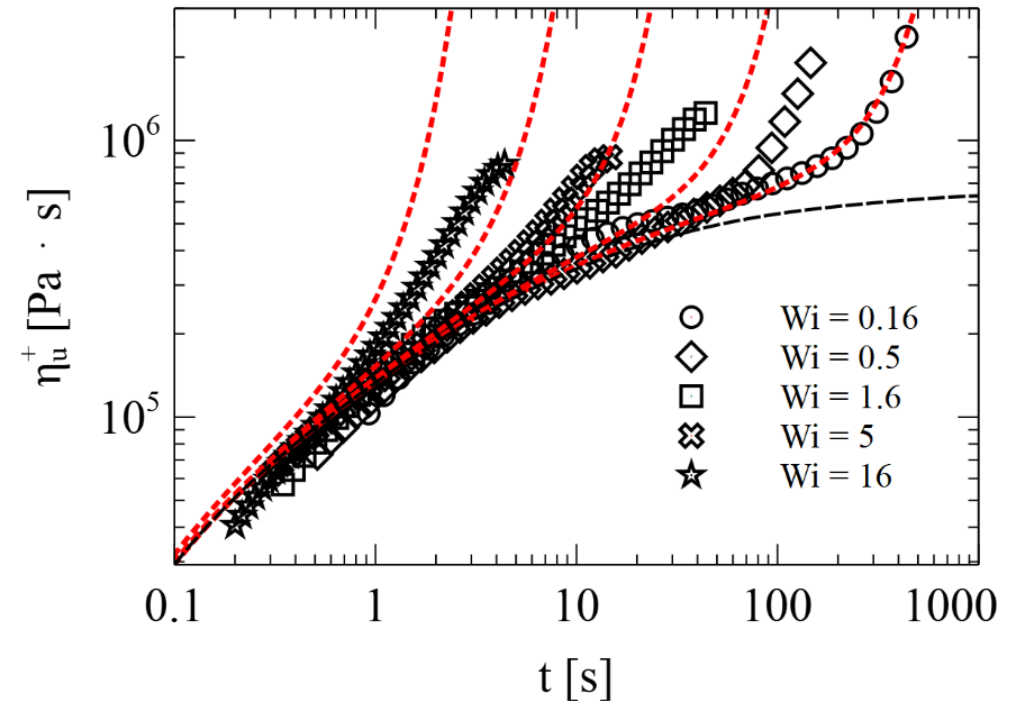
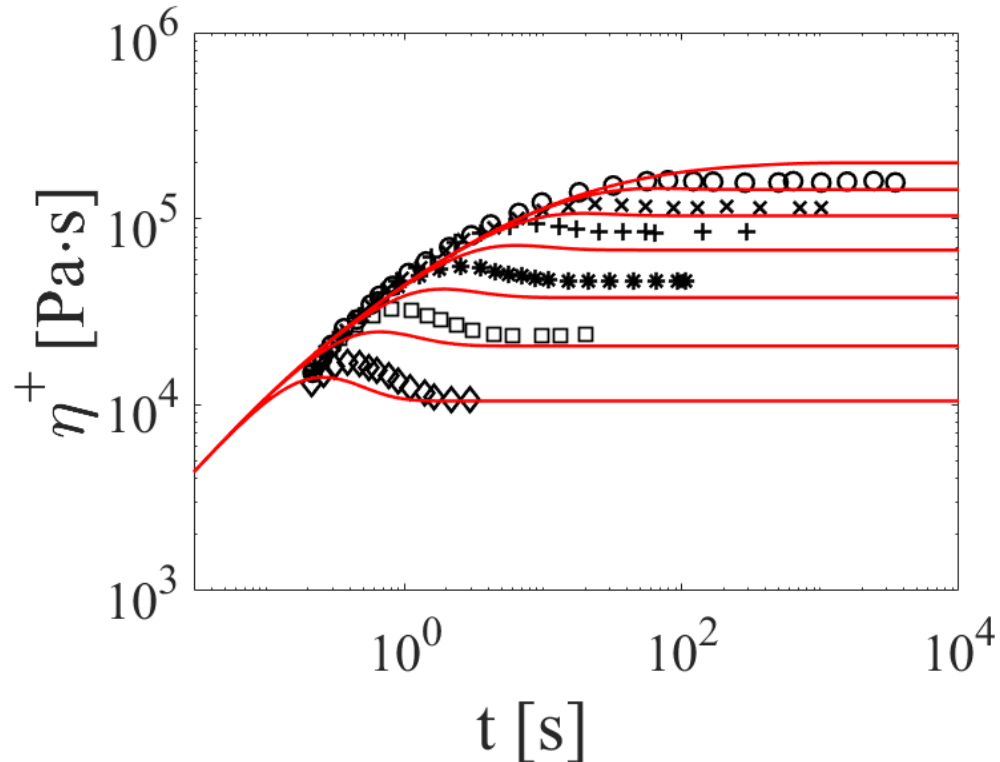
Constitutive Model: Rolie Poly

Graham et al. JOR 2003
Likhtman et al. JNNFM 2003

- Rolie Poly Model: ROuse Linear Entangled POLYmers

$$\frac{d\sigma}{dt} = \kappa \cdot \sigma + \sigma \cdot \kappa^T - \frac{1}{\tau_d} (\sigma - I) - \frac{2(1 - \sqrt{(3/\text{tr}\sigma)})}{\tau_R} \left(\sigma + \beta \left(\frac{\text{tr}\sigma}{3} \right)^\delta (\sigma - I) \right)$$

- Predictions



- Future Work: Implement Finite extensibility. Kabameni et al. Rheol Acta 2009

Thermal Hall Effect

1. Shear





2. Quench & Cut




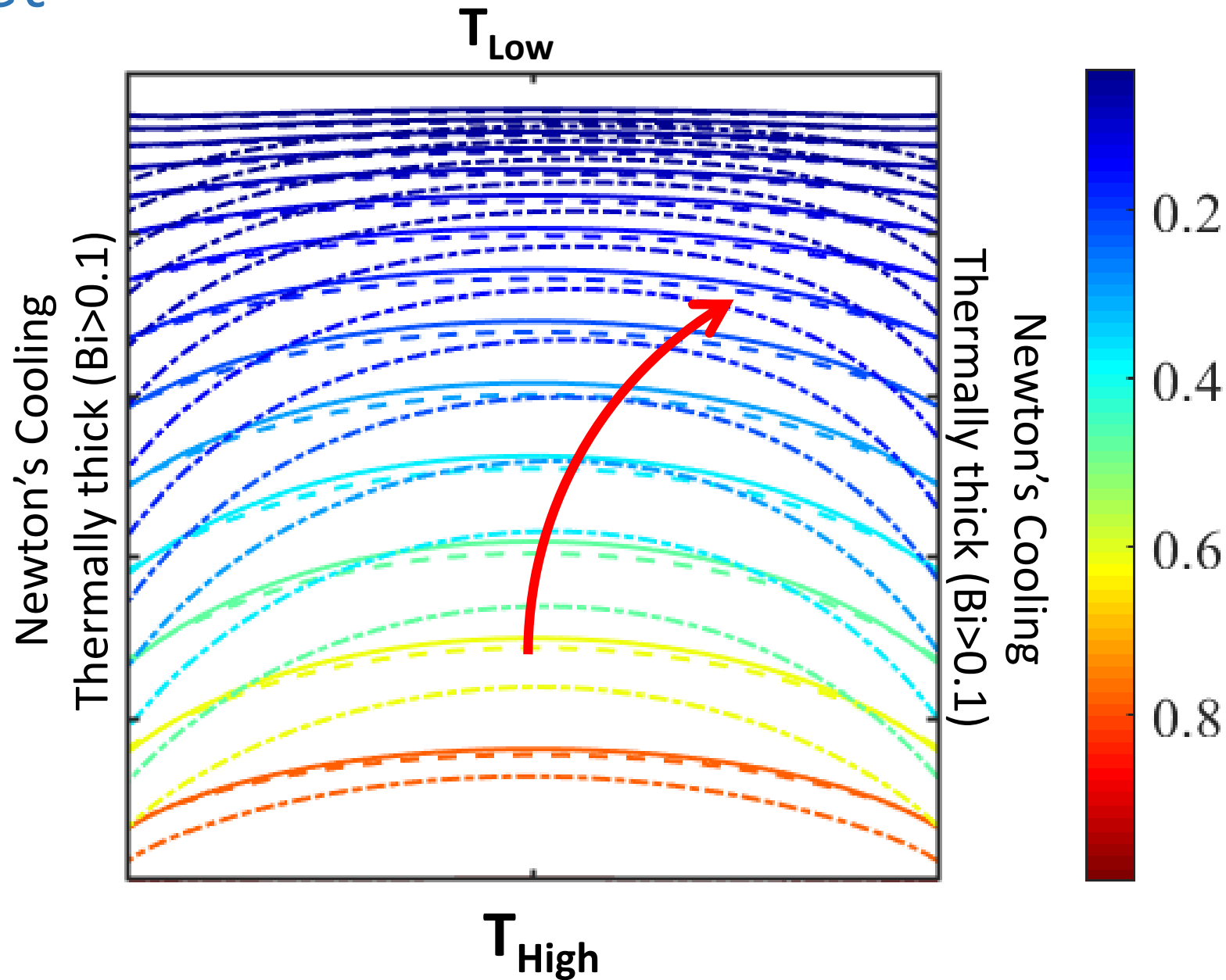
3. Subject to ∇T



 $\alpha_{11}=1.00, \alpha_{22}=1.00, \alpha_{12}=0.00$

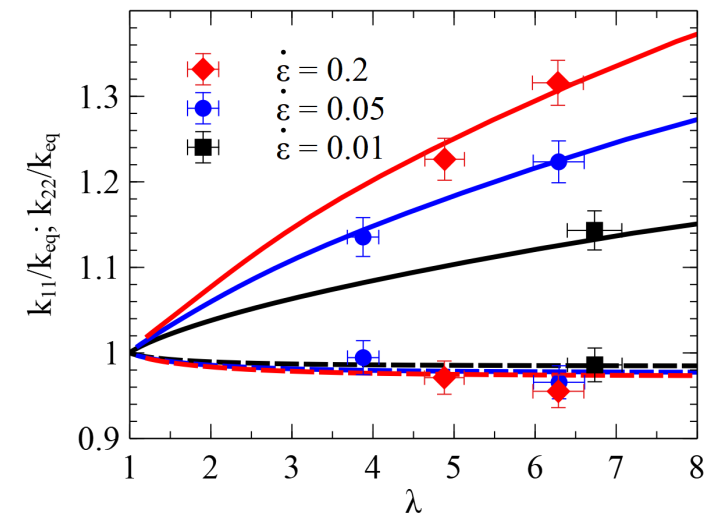
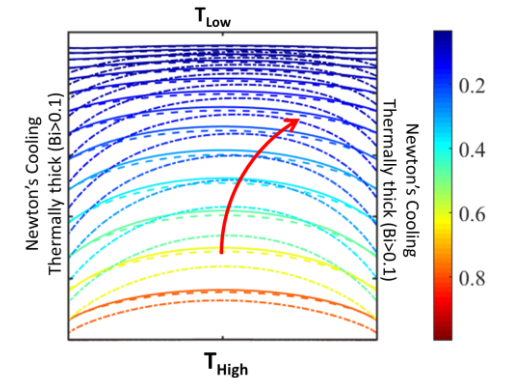
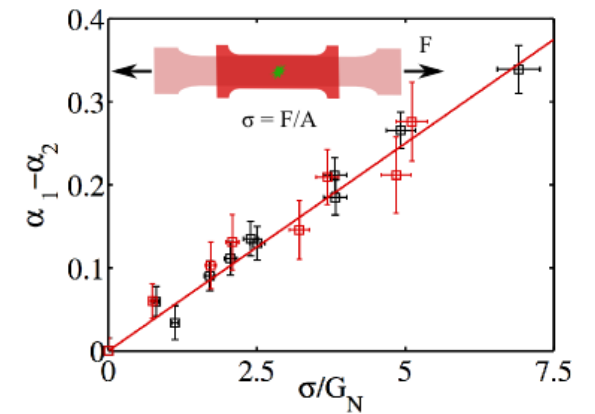
 $\alpha_{11}=1.20, \alpha_{22}=0.95, \alpha_{12}=0.00$

 $\alpha_{11}=1.20, \alpha_{22}=0.95, \alpha_{12}=0.25$



Conclusions

1. Thermal transport becomes anisotropic in polymers subjected to deformation
2. Flow induced anisotropy has significant implications in polymer processing
3. Experimental evidence of:
 - Proportionality to Stress: Stress-Thermal Rule (STR)
 - Universality
 - Beyond Finite Extensibility
4. We can use constitutive models (XPP, RP...), that are amenable to numerical flow simulations, in combination with the STR to include anisotropy in thermal conductivity in non-isothermal flows



Thank you!

David C. Venerus and Jay D. Schieber (Illinois Institute of Technology)
NSF Grant Nos. DMR-706582 and CBET-1336442.

Wilco M.H. Verbeeten (Universidad de Burgos)

Molecular to Continuum Investigation of Anisotropic Thermal Transport in Polymers
“MCIATTP”
Project # 750985



UNIVERSIDAD
DE BURGOS



MARIE CURIE ACTIONS



