

A Combination of the eXtended Pom-Pom Model and the Stress-Thermal Rule to Predict Anisotropy in Thermal Conductivity in Non-Linear Polymeric Flows

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MARIE CURIE **ACTIONS**



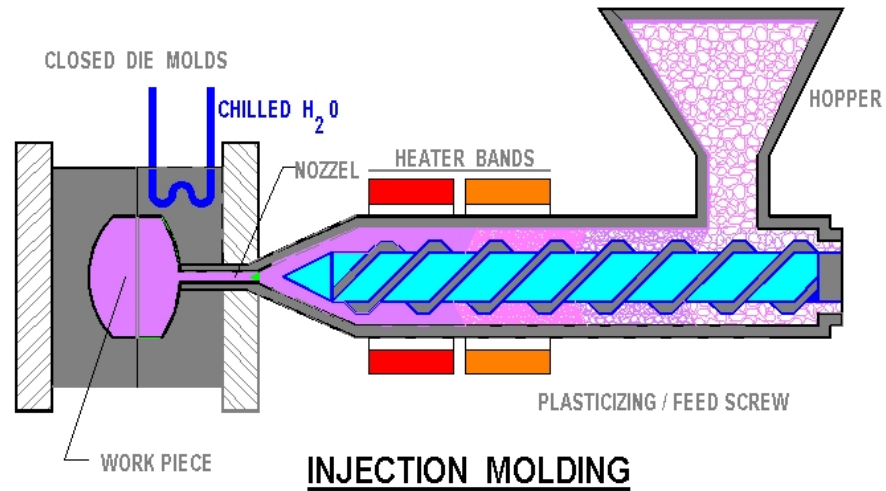
1. Thermal transport becomes anisotropic in polymers subjected to deformation
2. Experimental evidence of:
 - Proportionality to Stress: Stress-Thermal Rule (STR)
 - Universality
 - Beyond Finite Extensibility
3. eXtended Pom-Pom (XPP) as a constitutive model amenable for FEM simulation
4. Predictions of anisotropy for uniaxial and shear flows through combination of STR and XPP

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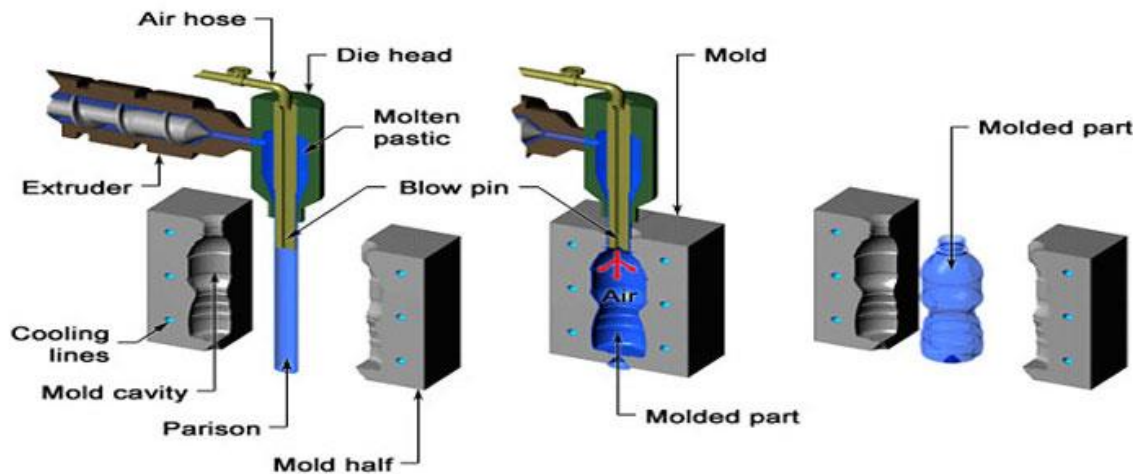
Motivation: Polymer Processing



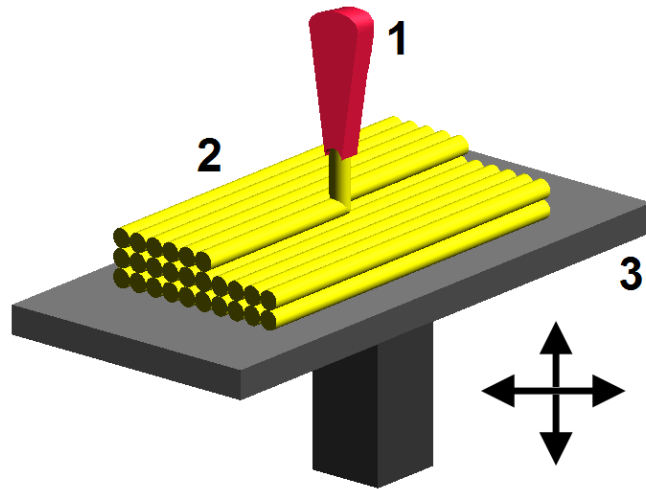
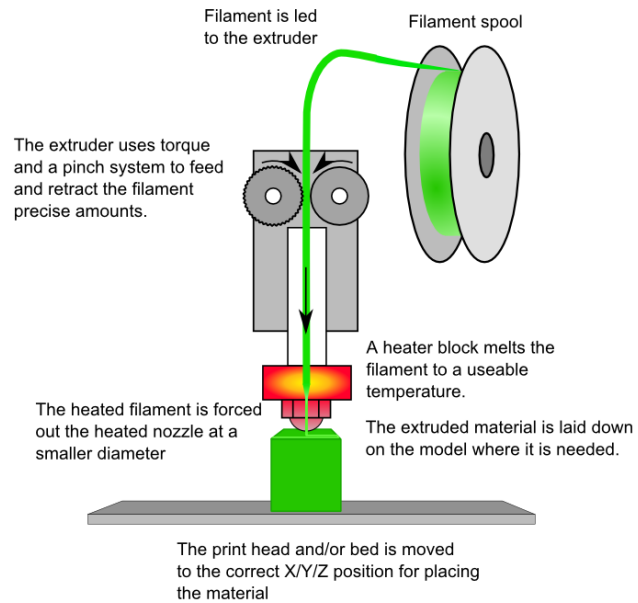
Global plastics market is expected to reach 654 billion USD by 2020

Thermal Transport Affects:

- Injection Pressure
- Cavity Flow
- Residual Stress
- Part Shrinkage



Motivation: Polymer Processing

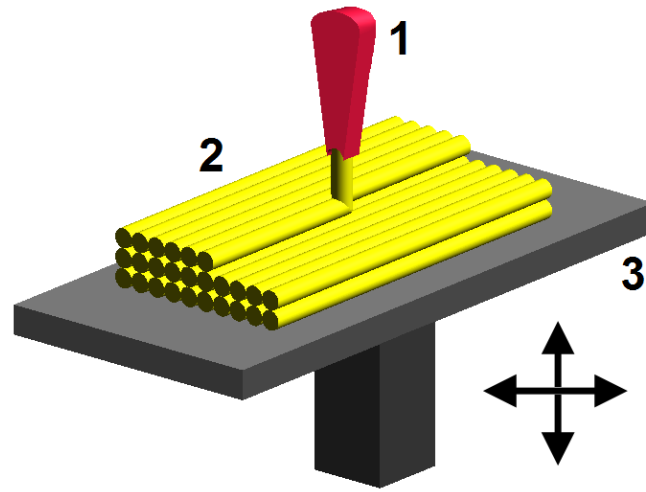
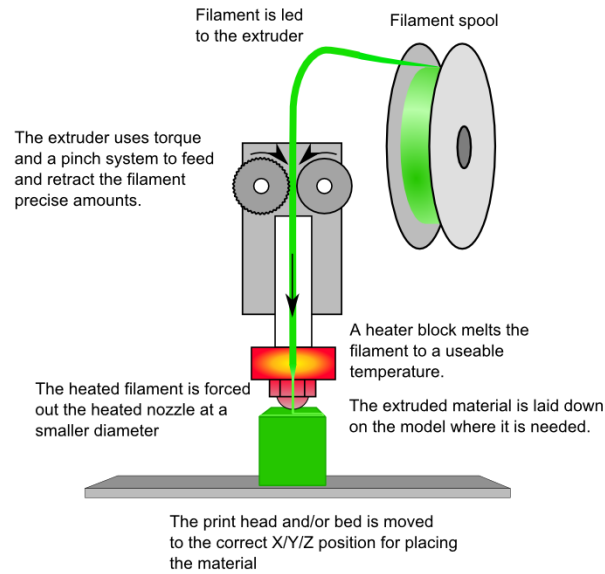


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Thermal Transport Affects:

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Develop a Molecular to Continuum methodology to better understand and simulate this kind of flows

Non-Isothermal Transport Phenomena

Balance Equations:

$$\text{Mass: } \frac{\partial \rho}{\partial t} = -\nabla \cdot (\rho \mathbf{v})$$

$$\text{Momentum: } \frac{\partial \rho \mathbf{v}}{\partial t} = -\nabla \cdot (\rho \mathbf{v} \mathbf{v} + \boldsymbol{\pi})$$

$$\text{Internal Energy: } \frac{\partial \rho \hat{u}}{\partial t} = -\nabla \cdot (\rho \hat{u} \mathbf{v} + \mathbf{q}) - \boldsymbol{\pi} : \nabla \mathbf{v}$$

Constitutive equations:

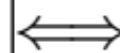
$$\mathbf{q} = -k \nabla T$$

$$\hat{c}_v = \hat{c}_v(T)$$

$$\boldsymbol{\tau} = \eta(T) [\nabla \mathbf{v} + \nabla \mathbf{v}^T]$$

- High stresses & Low thermal conductivity.

Mechanical behavior and flow



Thermal properties

Anisotropic Thermal Conduction

Fourier's Law: Thermal transport in deformed polymers is diffusive and anisotropic.

$$\mathbf{q} = -\mathbf{k} \cdot \nabla T$$

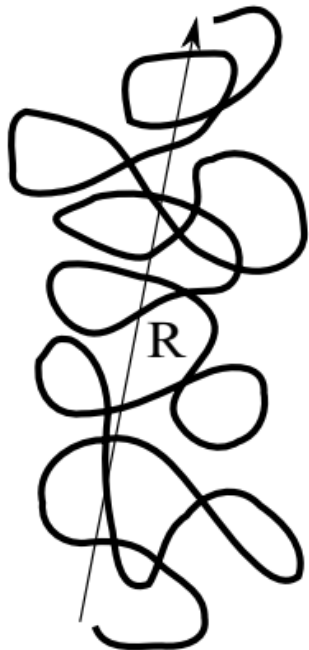
\mathbf{k} is a tensor!

Observation: k_{eq} increases with molecular weight.

Ueberreiter & Otto-Laupenmühlen, Kolloid Z. 1953

Hypothesis: *Energy transport along the backbone of a polymer chain is more efficient than between chains.*

Simple molecular arguments:



$$\mathbf{k} \propto \langle \mathbf{R}\mathbf{R} \rangle \quad + \quad \boldsymbol{\tau} \propto \langle \mathbf{R}\mathbf{R} \rangle$$

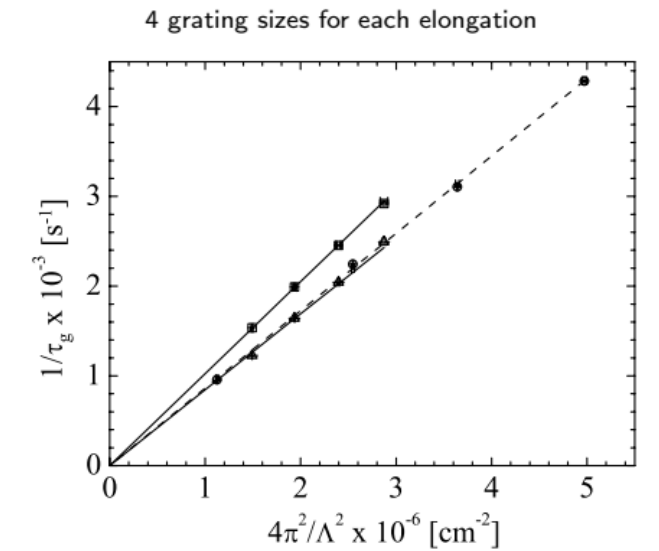
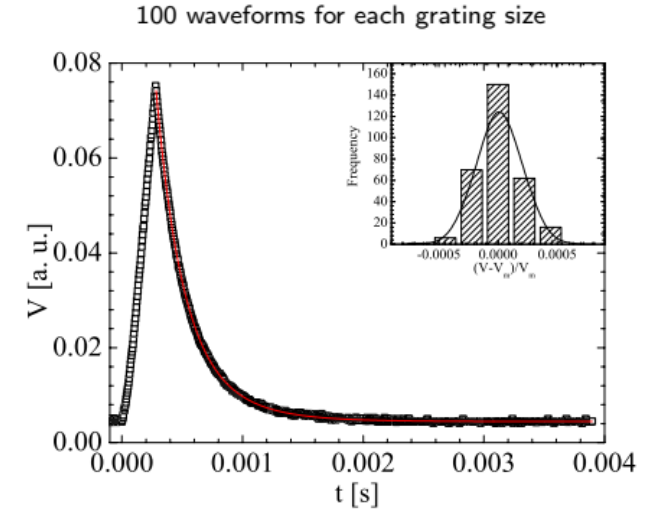
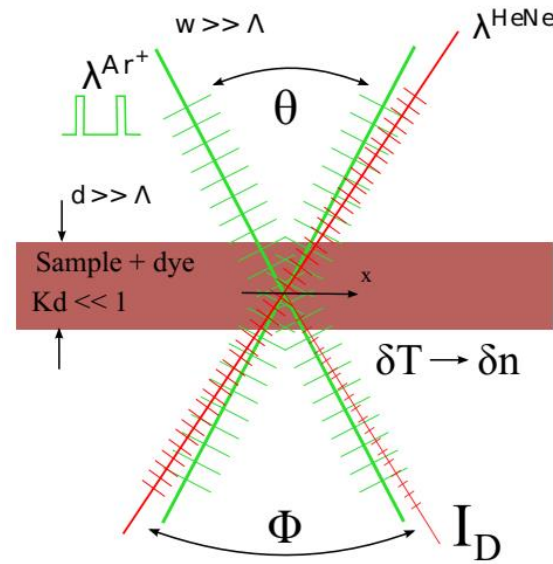
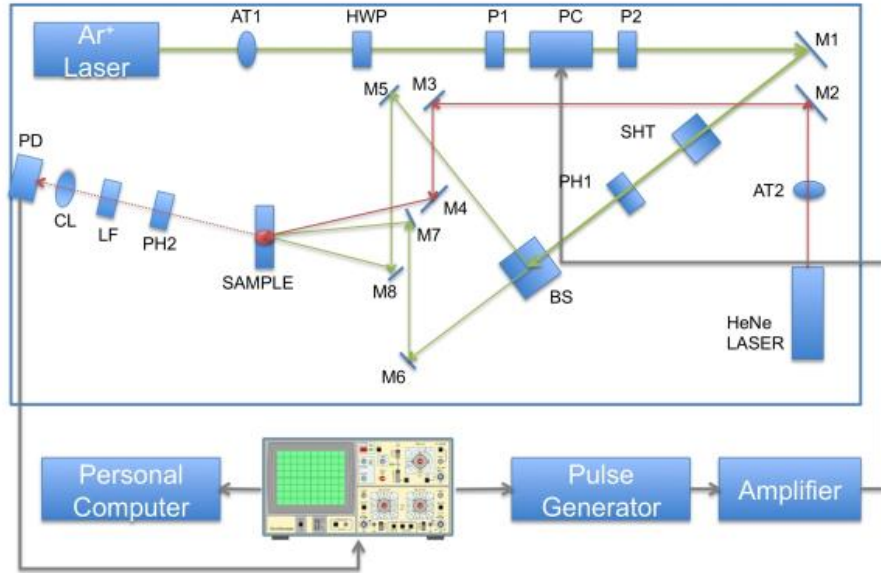
$$\mathbf{k} - \frac{1}{3}\text{tr}(\mathbf{k})\boldsymbol{\delta} = k_{\text{eq}}C_t \left[\boldsymbol{\tau} - \frac{1}{3}\text{tr}(\boldsymbol{\tau})\boldsymbol{\delta} \right]$$

The Stress-Thermal Rule

B.H.A.A. van den Brule, Rheol Acta 1989.
Öttinger and Petrillo, J. Rheol. 40 (5) 1996.
Curtiss and Bird, J. Chem. Phys. 107 (13) 1997.

$$C_t \propto \frac{nk_B^2 T}{\zeta}$$

Experiments: Forced Rayleigh Scattering (FRS)

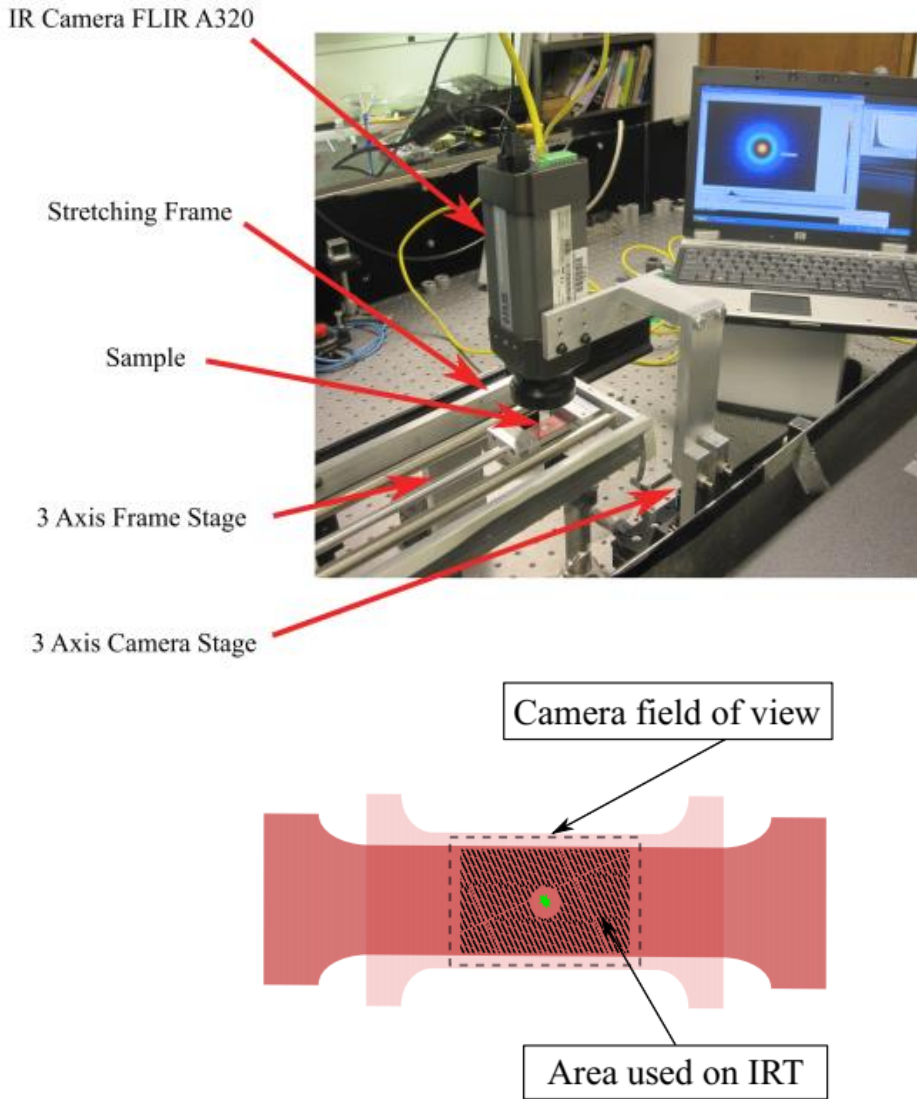


Intensity/Voltage at the photodetector:

$$V(t) = A \exp\left(-2\frac{t}{\tau_g}\right) + B \exp\left(-\frac{t}{\tau_g}\right) + C$$

$$\frac{1}{\tau_g} = D_{th} \frac{4\pi^2}{\Lambda^2} \quad D_{th} = \frac{k}{\rho \hat{c}_p}$$

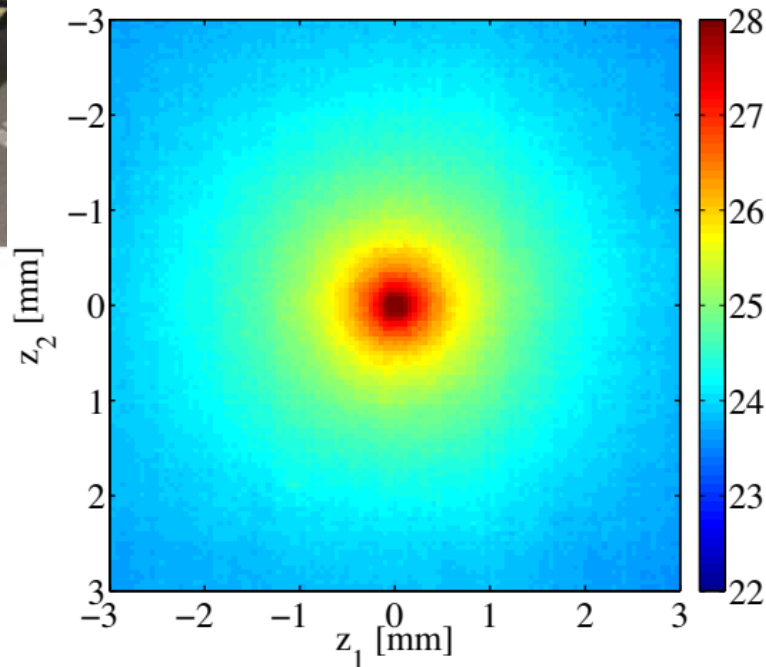
Experiments: Infrared Thermography (IRT)



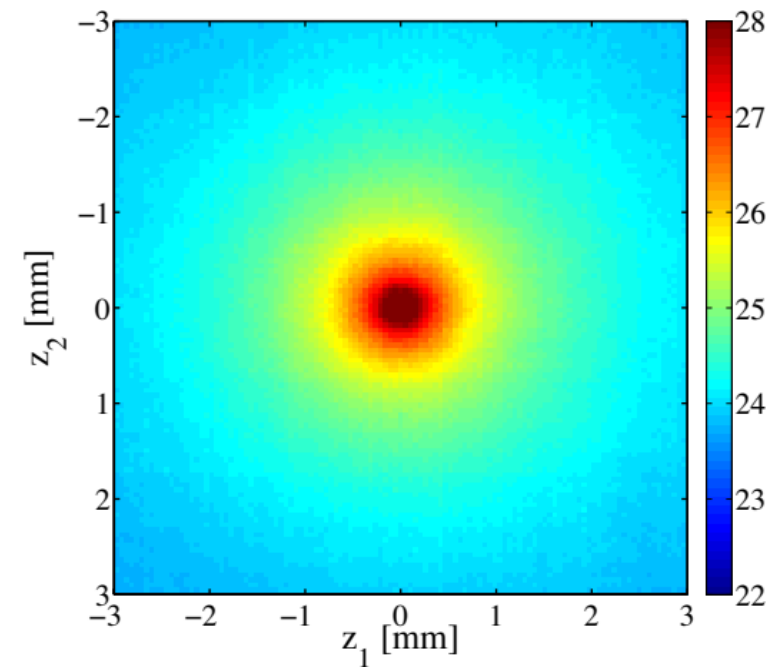
$$\theta(x_1, x_2) = \frac{1}{4\sqrt{\alpha_1\alpha_2}} K_0 \left(\sqrt{2\text{Bi}(x_1^2/\alpha_1 + x_2^2/\alpha_2)} \right)$$

$$KI_0w^2/k_{\text{eq}}, \quad \text{Bi} = hd/k_{\text{eq}}$$

$$\alpha_1 = k_{11}/k_{\text{eq}}, \quad \alpha_2 = k_{22}/k_{\text{eq}}$$

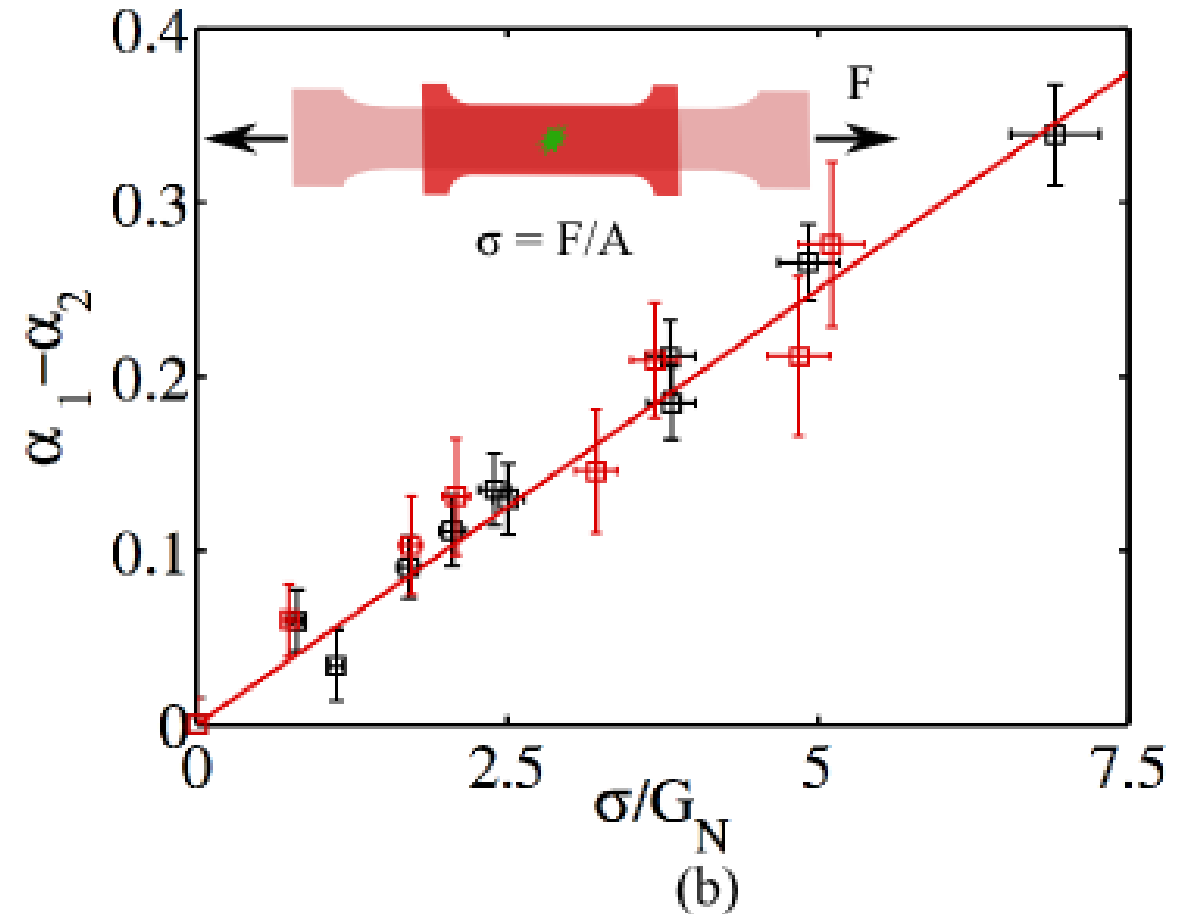
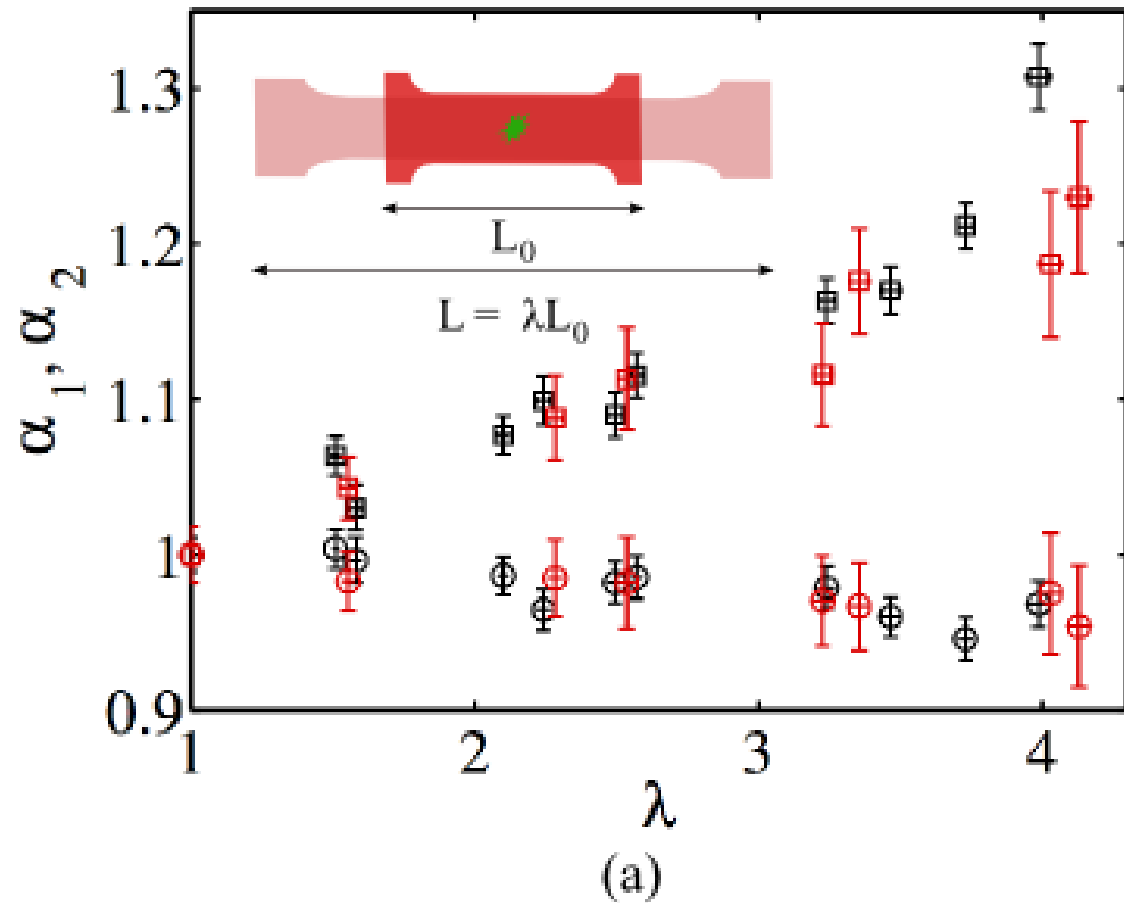


Un-stretched sample, $\lambda = 1$
and $\text{Bi}_0 = 0.029 \pm 0.001$



Stretched sample, $\lambda = 4.129$,
 $\alpha_1 = 1.23 \pm 0.049$ and $\alpha_2 = 0.954 \pm 0.039$

Comparison FRS and IRT



Key Findings: Universality...

Stress-Thermal Coefficients for several polymeric materials

Material	Deformation –	G_N [kPa]	$C_t \times 10^4$ [kPa ⁻¹]	$C_t G_N$ –	$C \times 10^9$ [Pa ⁻¹]
PIB 85k ⁷	Shear	320 ¹	1.9	0.061 ± 0.024	1.45
PIB 130k ⁷	Shear	320 ¹	1.2	0.038 ± 0.022	1.45
xI-PDMS ⁶	Uniax.	200 ¹	1.3	0.026 ± 0.008	0.13-0.26
xI-PBD 200k ⁵	Uniax.	760 ¹	0.73	0.051 ± 0.011	3.5
xI-PBD 150k ⁵	Uniax.	760 ¹	0.93	0.059 ± 0.014	3.5
xI-PI 100k ⁴	Uniax.	370 ²	0.37	0.014 ± 0.005	2.2
PS 260k ³	Uniax.	200 ¹	1.65	0.033 ± 0.007	-4.8
PMMA 83k ³	Uniax.	310 ¹	1.7	0.054 ± 0.011	0.16

$$C_t G_N \sim 0.04$$

- (1) Fetters et al. Macromolecules 27, 17 (1994)
- (2) Fetters et al. Macromolecules 37 (2004)
- (3) Gupta et al. Journal of Rheology 57 (2013)
- (4) Nieto Simavilla et al. J. Pol. Sci. B 50 (2012)
- (5) Venerus et al. Macromolecules 42 (2009)
- (6) Broerman et al. J.Chem. Phys. 111 (1999)
- (7) Venerus et al. Phys. Rev. Lett. 82 (1999)

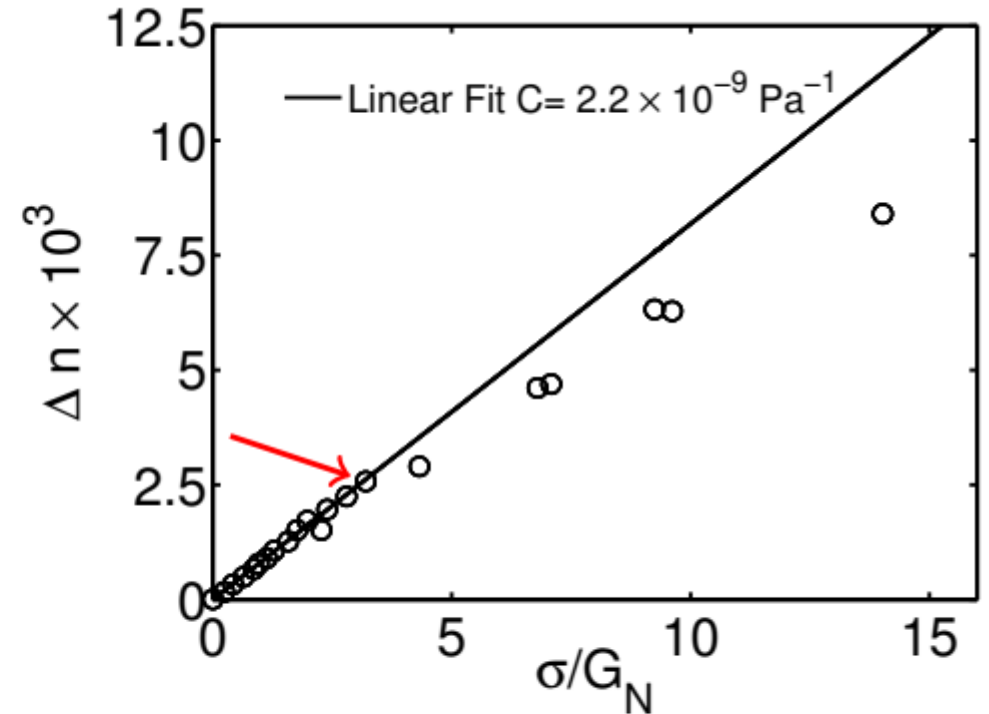
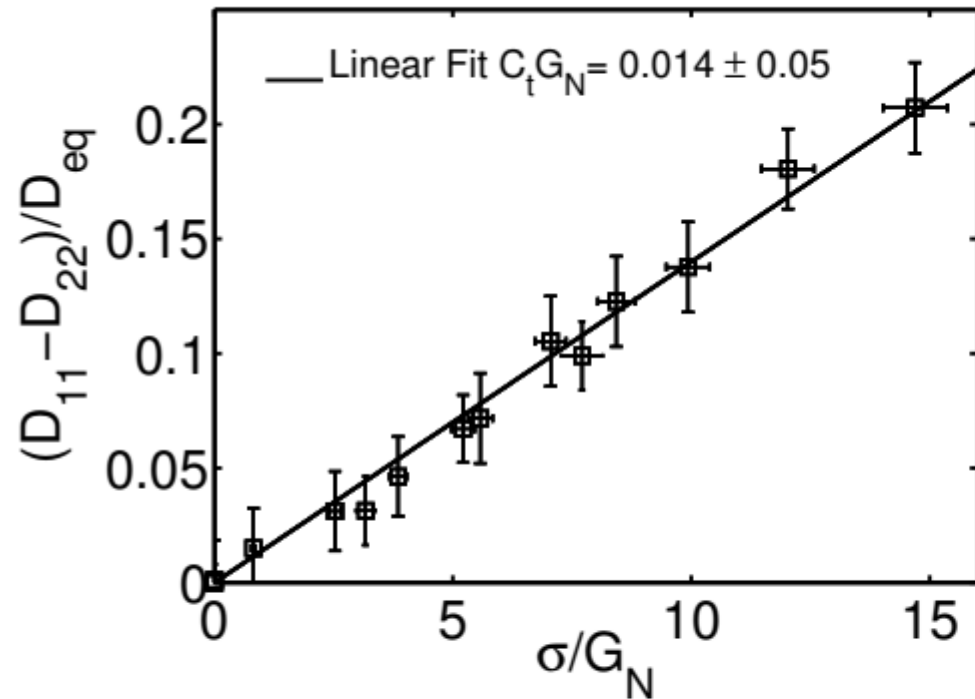
Stress-thermal Rule:

$$\mathbf{k} - \frac{1}{3}\text{tr}(\mathbf{k})\boldsymbol{\delta} = k_{\text{eq}}C_t(\boldsymbol{\tau} - \frac{1}{3}\text{tr}(\boldsymbol{\tau})\boldsymbol{\delta})$$

Stress-optic Rule:

$$\mathbf{n} - \frac{1}{3}\text{tr}(\mathbf{n})\boldsymbol{\delta} = C(\boldsymbol{\tau} - \frac{1}{3}\text{tr}(\boldsymbol{\tau})\boldsymbol{\delta})$$

Key Findings: ...Beyond Finite Extensibility



The STR stays valid where the SOR fails!

Constitutive Model: eXtended Pom-Pom

- What physics are in the model?

$$\overset{\nabla}{\boldsymbol{\tau}} + \boldsymbol{\lambda}(\boldsymbol{\tau})^{-1} \cdot \boldsymbol{\tau} - 2G_0 \mathbf{D}_u = \mathbf{0}$$

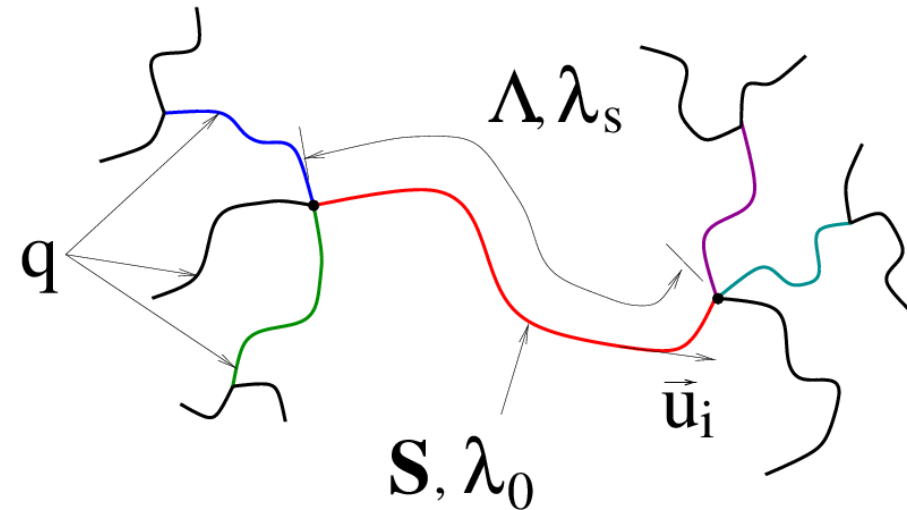
$$\alpha \neq 0 \rightarrow \Psi_2 \neq 0$$

$$\boldsymbol{\lambda}(\boldsymbol{\tau})^{-1} = \frac{1}{\lambda_{0b}} \left[\frac{\alpha}{G_0} \boldsymbol{\tau} + f(\boldsymbol{\tau})^{-1} \mathbf{I} + G_0 (f(\boldsymbol{\tau})^{-1} - 1) \boldsymbol{\tau}^{-1} \right] \quad \Lambda = \sqrt{1 + \frac{I_{\boldsymbol{\tau}}}{3G_0}}$$

$$\frac{1}{\lambda_{0b}} f(\boldsymbol{\tau})^{-1} = \frac{2}{\lambda_s} \left(1 - \frac{1}{\Lambda}\right) + \frac{1}{\lambda_{0b}} \left(\frac{1}{\Lambda^2} - \frac{\alpha I_{\boldsymbol{\tau} \cdot \boldsymbol{\tau}}}{3G_0^2 \Lambda^2} \right) \quad \lambda_s = \lambda_{0s} e^{-\frac{2}{q}(\Lambda-1)}$$

- Why XPP?

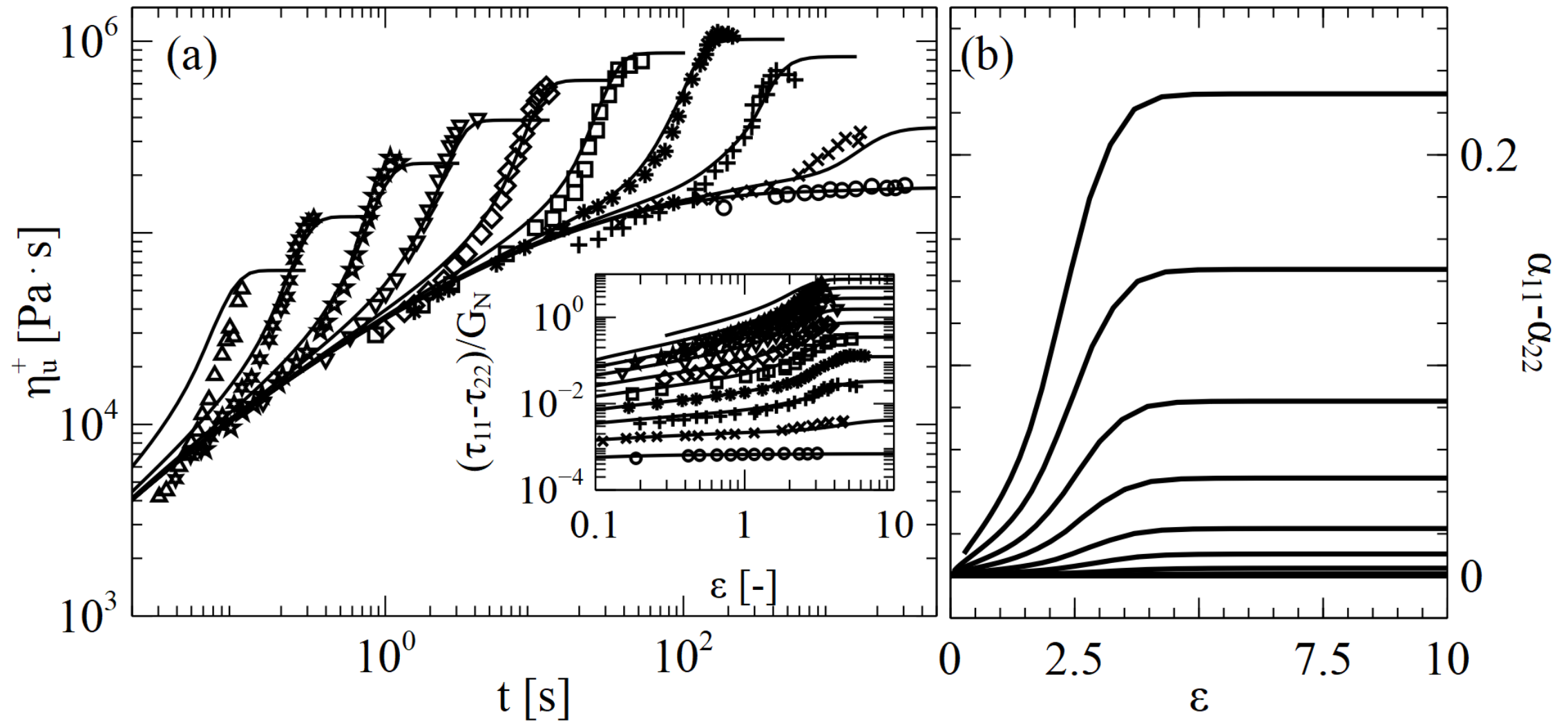
- Amenable to FEM
- Able to describe non-linear rheology
- X: Avoids finite extensibility discontinuities
- X: Includes second normal stress difference



Data: IUPAC_A LDPE melt at 170°C

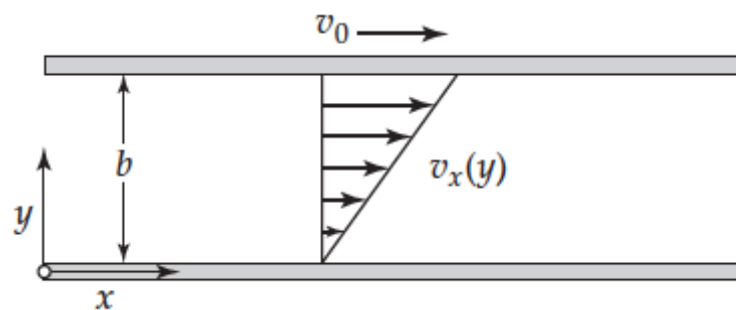
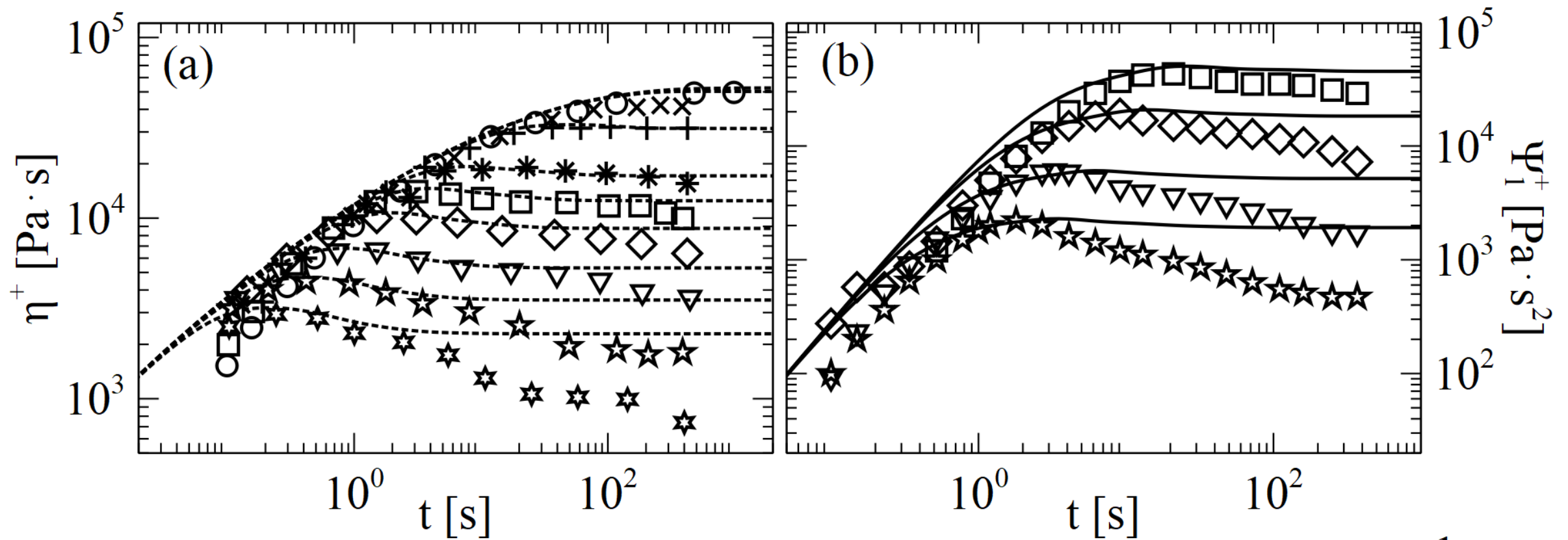
PP: McLeish and Larson. JOR 1998
xPP: Verbeeten et al. JOR 2001

Transient Start-up: Uniaxial

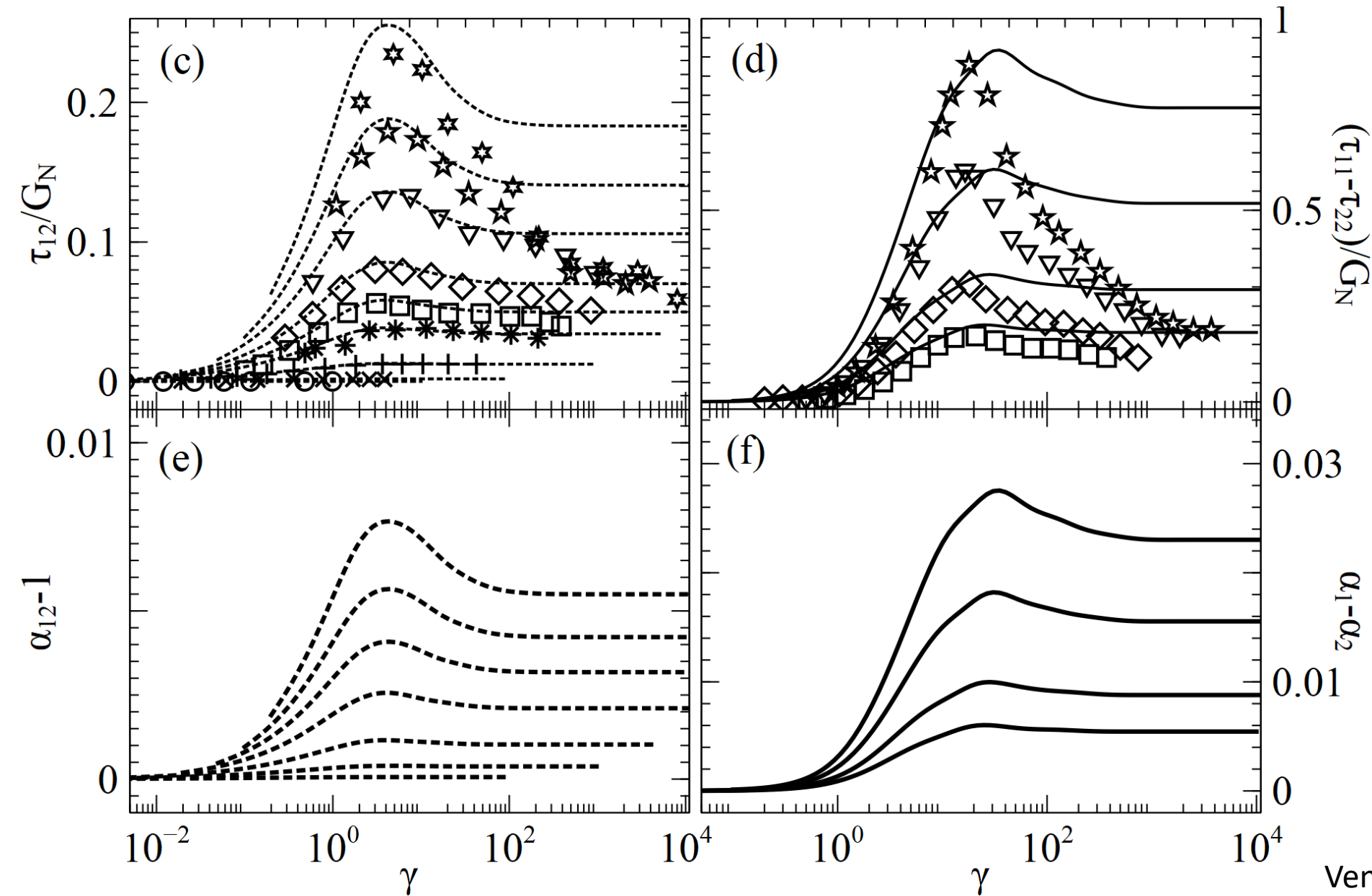


The anisotropy in TC is comparable to that observed in PS and PMMA melts $\sim 20\%$.
Gupta et al. Journal of Rheology 57, 2013.

Transient Start-up: Shear Rheology



Transient Start-up: Shear

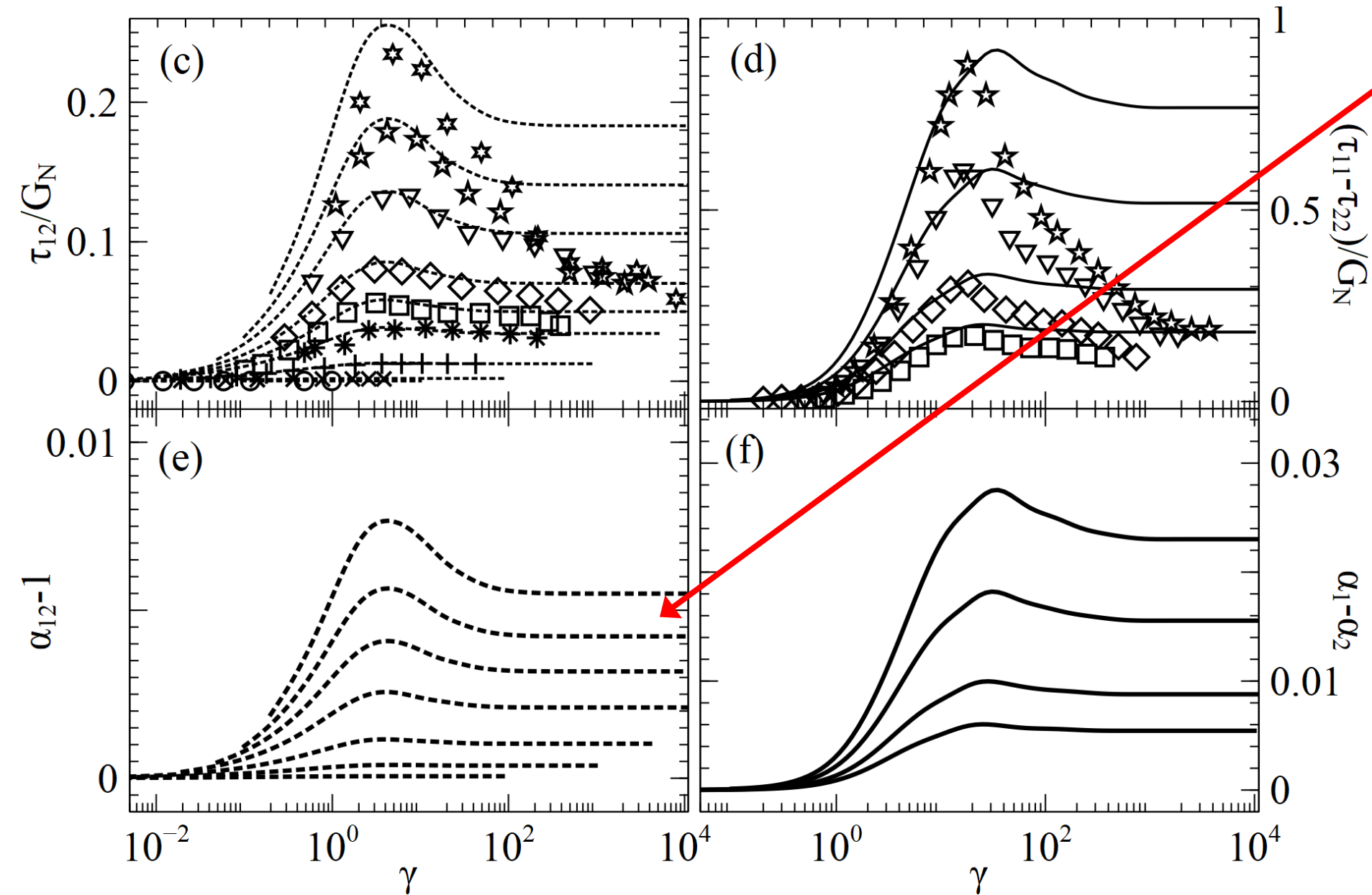


There is a non-zero off-diagonal component in shear flows

$$\mathbf{k} = \begin{pmatrix} k_{11} & k_{12} & k_{13} \\ & k_{22} & k_{23} \\ & & k_{33} \end{pmatrix}$$

$$\mathbf{q} = \mathbf{k} \cdot \nabla T$$

Transient Start-up: Shear



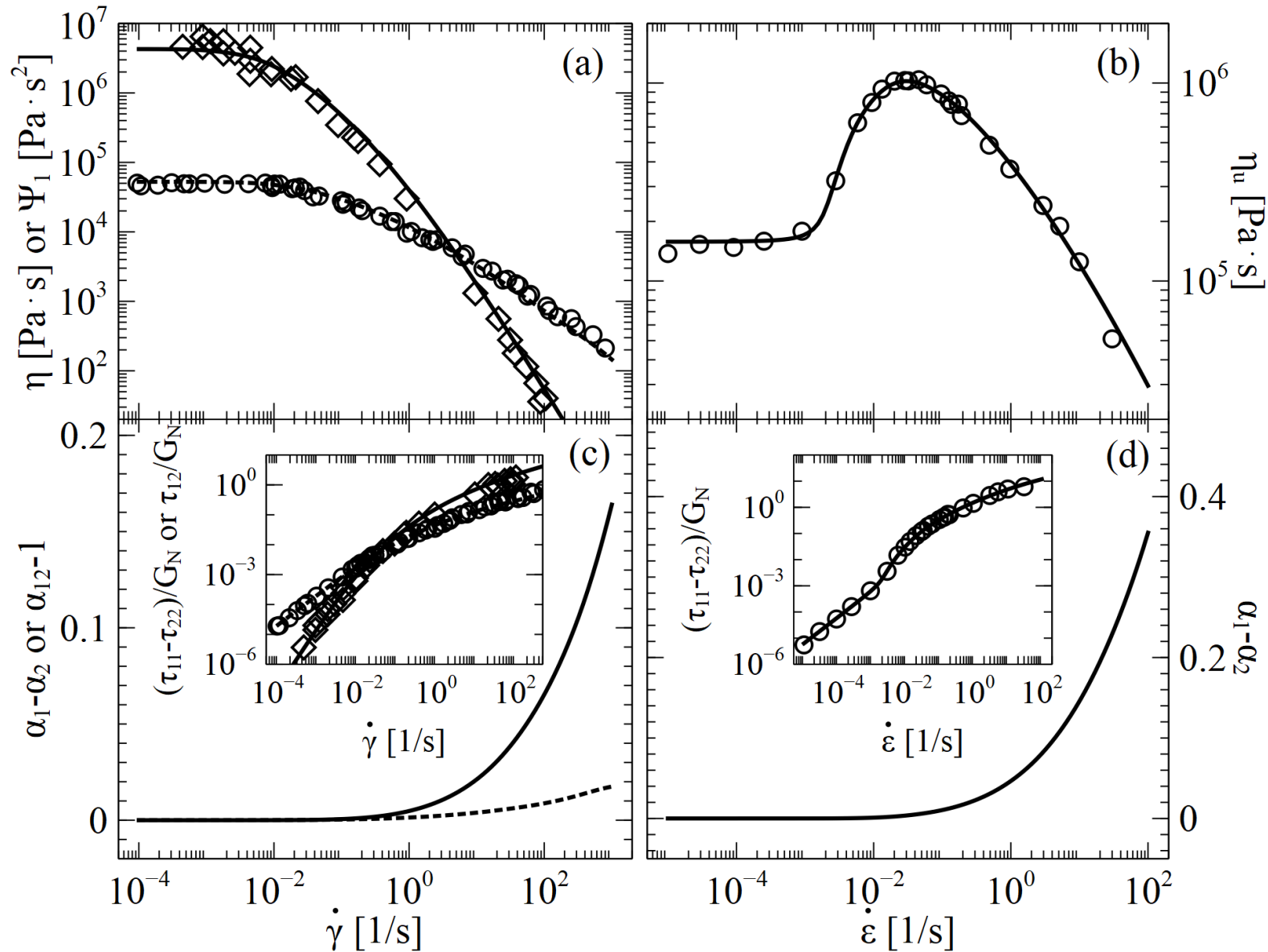
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A temperature gradient in the 1-direction can generate heat flow in the 2-direction:
Thermal Hall Effect

Steady-State: Shear and Uniaxial Extension



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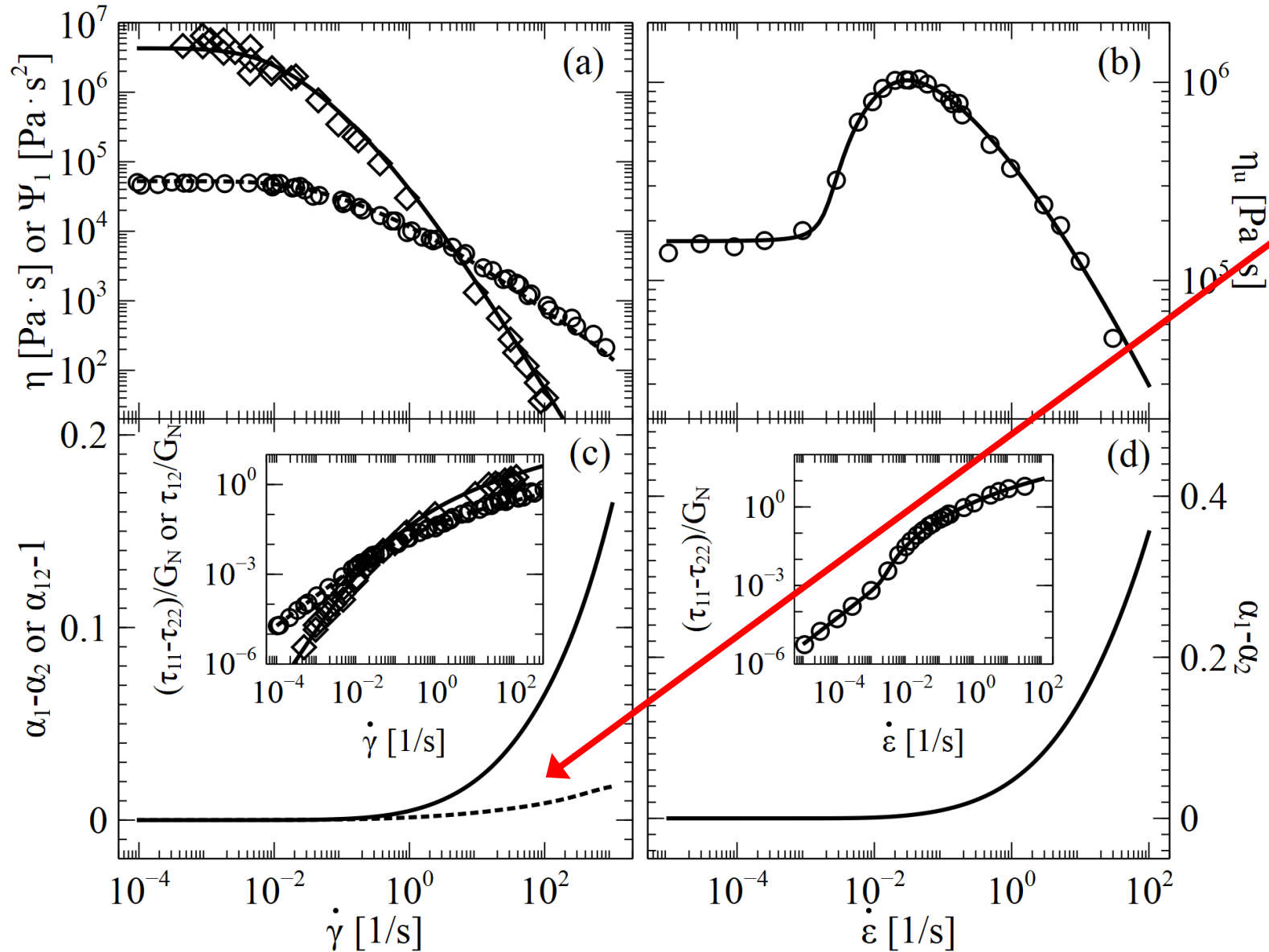
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Future Work: MTCIATTP project

A. Implementation of New Constitutive Models

- Other polymer architectures: Rolie Poly
- Compare predictions with available experimental data PS, PMMA...

B. Develop a deeper understanding at the molecular level

- Why universal?
- Why beyond finite extensibility?

C. Implementation of Non-Isothermal Flow Simulations (FEM)

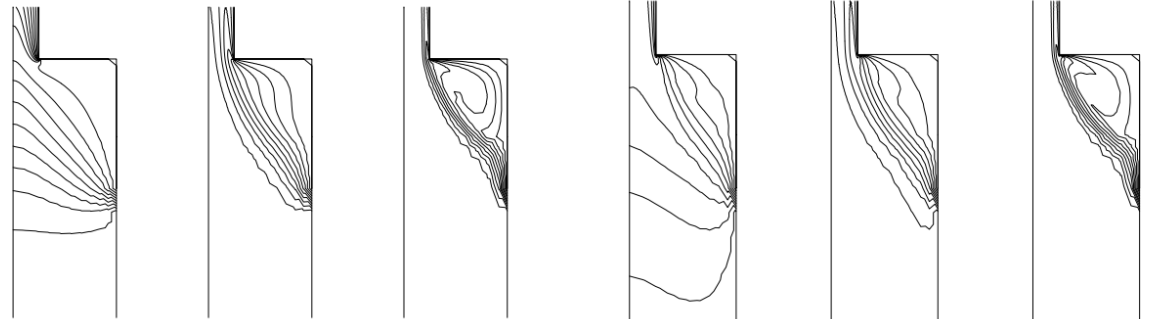


Isotropic Thermal Conductivity: k

Anisotropic Thermal Conductivity: \mathbf{k}

$$\mathbf{q} = -k\nabla T$$

$$\mathbf{q} = -\mathbf{k} \cdot \nabla T$$



Isotherms for Pe = 10, Pe = 100 and Pe = 1000.

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Conclusions

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Thank you!

David C. Venerus and Jay D. Schieber (Illinois Institute of Technology)
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Wilco M.H. Verbeeten (Universidad de Burgos)

Molecular to Continuum Investigation of Anisotropic Thermal Transport in Polymers
“MCIATTP”
Project # 750985



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